The background of the slide is a grayscale aerial photograph of a university campus. The image shows a dense cluster of modern buildings with various architectural styles, including some with glass facades and others with more traditional brickwork. There are several large parking lots filled with cars, and green lawns with trees and walkways in between the buildings. The overall scene is a typical college campus.

# MHD with embedded particle-in-cell model and its applications

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# Outline

## ➤ Introduction

- Background of developing MHD with embedded particle-in-cell (MHD-EPIC) model

## ➤ Coupling between MHD and PIC

## ➤ Improvement of the PIC component

- Gauss's Law satisfying Energy Conserving Semi-Implicit Method (GL-ECSIM)
- FFlexible Exascale Kinetic Simulator (FLEKS)

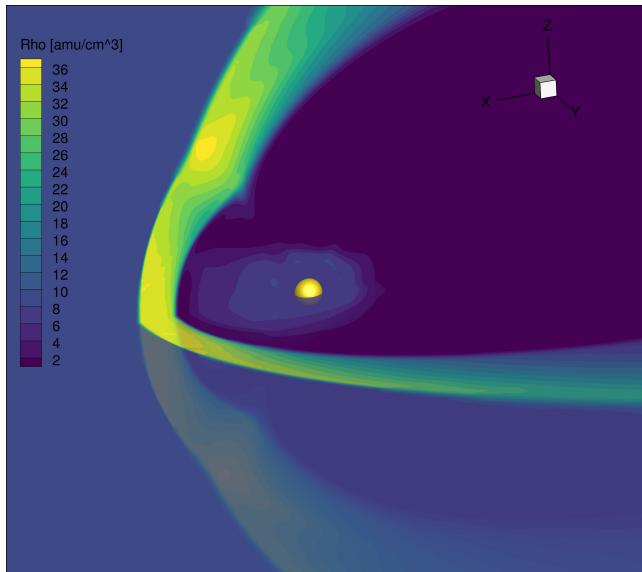
## ➤ Applications

- Mercury's magnetotail
- Earth's dayside magnetopause

# Introduction: MHD vs. Kinetic

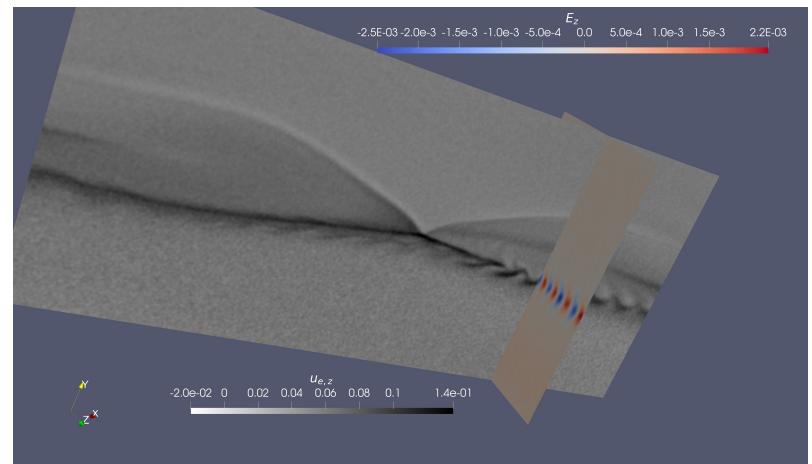
## ➤ Ideal MHD

- Efficient for 3D global simulations.
- Does not contain kinetic physics.



## ➤ Kinetic models of collisionless plasma: Vlasov-Maxwell solver, PIC

- Solve the evolution of the velocity space (Vlasov equation)
- Computationally expensive



## ➤ Why it is difficult to capture kinetic processes in global simulations?

- Orders of difference between the kinetic scales and the global scale
- Example of Earth's magnetosphere:

$\lambda_D \sim 100$  m;  $d_e \sim 5$  km;  $d_i \sim 100$  km; Magnetosphere size  $\sim 10R_E \sim 10^5$  km

# Introduction: Incorporate Kinetic Physics into Global Simulations

## ➤ Hybrid models

- Fluid massless electrons + kinetic (PIC or Vlasov solver) ions
- Local kinetic + global fluid -> **MHD-EPIC**

## ➤ Extended MHD

- Five-moment (isotropic pressure), six-moment (anisotropic pressures), and ten-moment (pressure tensors)

Earth's magnetosphere:

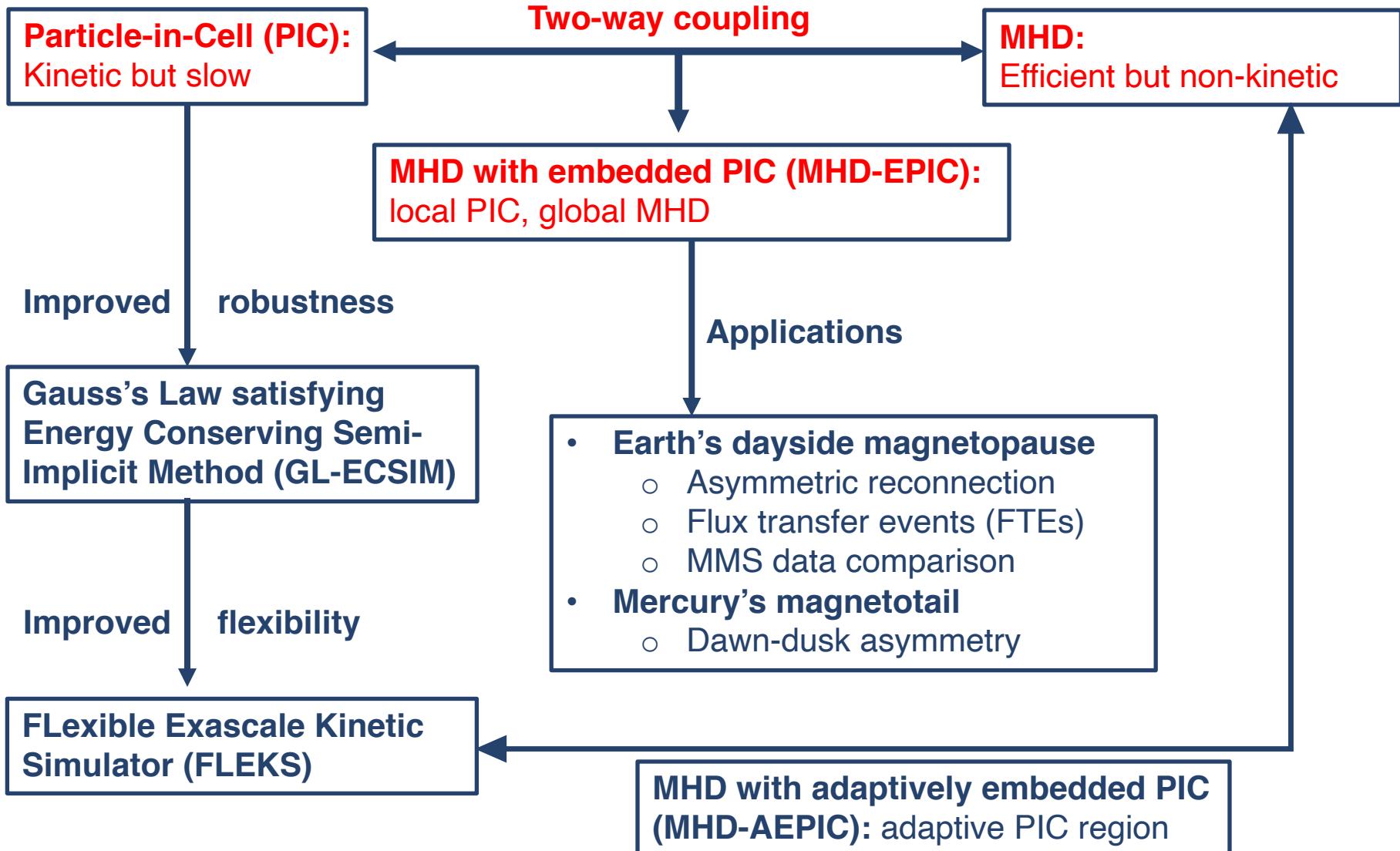
$\lambda_D \sim 100$  m;  $d_e \sim 5$  km;  $d_i \sim 100$  km;

Magnetosphere size  $\sim 10R_E \sim 10^5$  km

## ➤ Multi-scale simulations: reduce the separation between scales

- Reduce speed of light: reduce separation between  $\lambda_D$  and  $d_e$
- Reduce the  $m_i/m_e$  ratio: reduce separation between  $d_e$  and  $d_i$

# Outline



# MHD-EPIC Algorithm

**The goal is to combine the efficiency of the global fluid code with the physics capabilities of the local PIC code!**

➤ **MHD code - BATS-R-US**

- Solves the ideal, Hall, two-fluid, multi-ion, anisotropic MHD equations on 1, 2 or 3D block-adaptive Cartesian and non-Cartesian grids
- Five- and six-moment multi-fluid equations
- Explicit and implicit time stepping, up to 5<sup>th</sup> order accurate schemes
- Complicated initial and boundary conditions, variety of source terms

➤ **PIC code - IPIC3D & FLEKS**

- Solves full set of Maxwell's Equations implicitly on 2D or 3D Cartesian grid
- Relaxed stability constraints compared to explicit PIC

➤ **Framework – SWMF**

- Executes and couples multiple models on massively parallel machines
- Efficient parallel coupler developed for MHD-EPIC

# Particle-in-Cell Method

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi \mathbf{J}$$

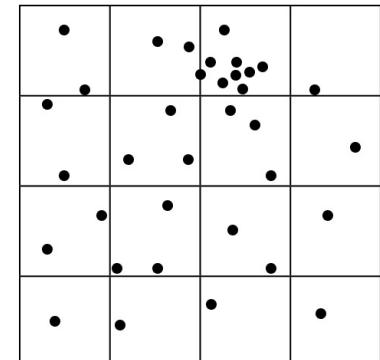
$$\frac{d\mathbf{p}}{dt} = q(\mathbf{E} + \frac{\mathbf{V} \times \mathbf{B}}{c})$$

## ➤ Physics

- Collisionless plasma
- Solves Maxwell's Equations
- Equations of motion for electrons and ion macro-particles

## ➤ Calculation steps

- Loop through all the particles once to calculate the current at each grid cell
- Solve the Maxwell equations on the grid
- Interpolate the EM field from the grid to the particles
- Move the particles



# Explicit PIC vs. Semi-Implicit PIC

## ➤ Explicit PIC

- Updates EM fields with known E, B and current
- Cell size needs to resolve Debye length to suppress finite grid instability
- Conditional stable: time step needs to resolve plasma oscillation and propagation of light wave

## ➤ Semi-Implicit PIC

- $0.5 \leq \theta \leq 1$
- Updates fields with E, B and J in the future.
- Semi-implicit:  $J^{n+1/2}$  linearly depends on  $E^{n+\theta}$ .  $J^{n+1/2} \sim J^n + M \cdot E^{n+\theta}$
- Solves linear system with an iterative solver
- No need to resolve Debye length
- Linearly unconditional stable: time step is not limited by plasma oscillation and propagation of light wave

$$\frac{\mathbf{B}^{n+1} - \mathbf{B}^n}{\Delta t} = -c\nabla \times \mathbf{E}^n$$

$$\frac{\mathbf{E}^{n+1} - \mathbf{E}^n}{\Delta t} = c\nabla \times \mathbf{B}^n - 4\pi \mathbf{J}^n$$

$$\frac{\mathbf{B}^{n+1} - \mathbf{B}^n}{\Delta t} = -c\nabla \times \mathbf{E}^{n+\theta}$$

$$\frac{\mathbf{E}^{n+1} - \mathbf{E}^n}{\Delta t} = c\nabla \times \mathbf{B}^{n+\theta} - 4\pi \mathbf{J}^{n+1/2}$$

# Explicit PIC vs. Semi-Implicit PIC

## ➤ Explicit PIC vs. Semi-implicit PIC

- Both solve the same set of equations, but with different discretization
- Converge to the same solution
- With the same cell size and number of macro-particles, explicit PIC is  $\sim 10$  times faster than semi-implicit PIC
- Explicit PIC needs to resolve Debye length and plasma oscillations for stability reason
- Semi-implicit PIC only needs to resolve the scales of interest

## ➤ Unresolved scales

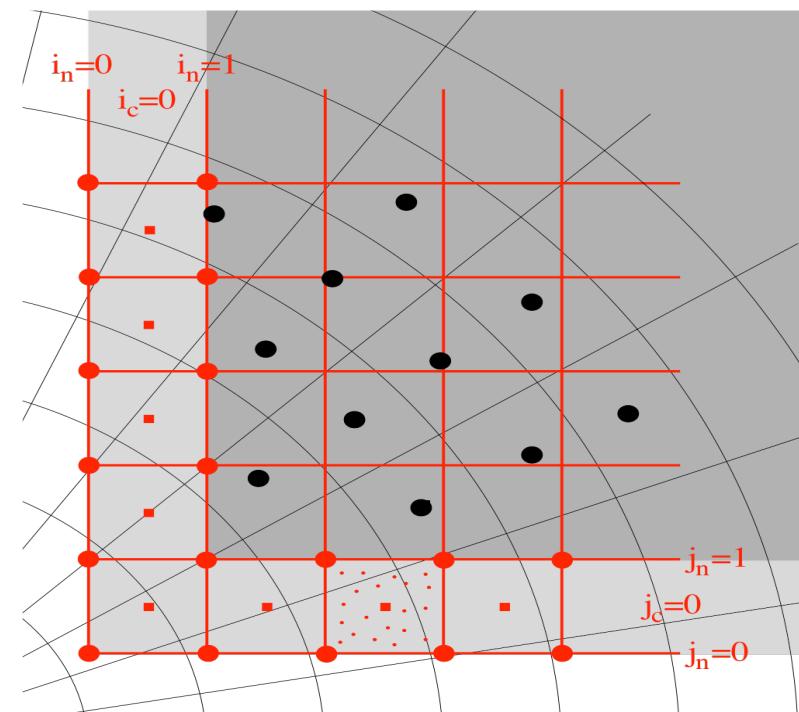
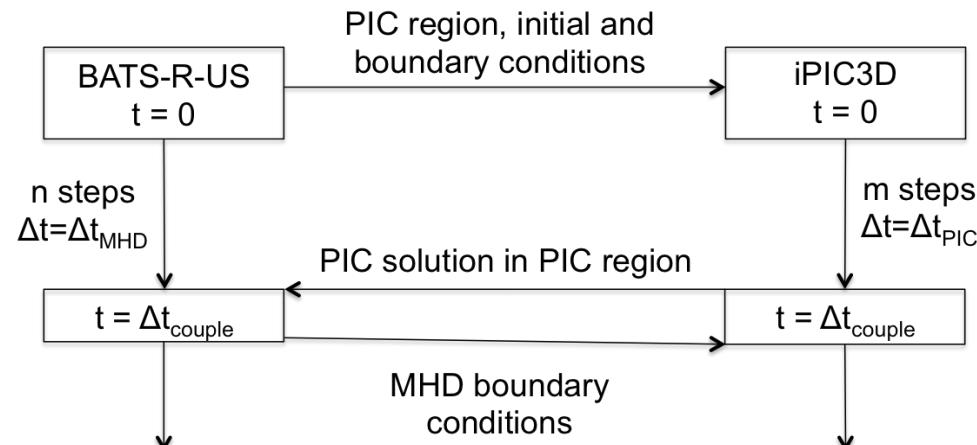
- When  $\Delta t$  is large, the unresolved waves are strongly damped (Brackbill et al., 1982, JCP)

# Coupling between MHD and PIC

## ➤ Capabilities

- Ideal, Hall and multi-ion MHD. Separate electron pressure.
- Different PIC and MHD grids including non-Cartesian AMR MHD grids.
- Different PIC and MHD time steps.
- Efficient coupling through the SWMF.
- Multiple PIC domains.
- Works in 2D and 3D.

➤ For the PIC code, MHD-EPIC is a method of setting initial and boundary conditions



# PIC Initialization

MHD equations:

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) &= 0 \\ \frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot \left[ \rho \mathbf{u} \mathbf{u} + p_{\perp} \mathbf{I} + (p_{\parallel} - p_{\perp}) \mathbf{b} \mathbf{b} - \frac{1}{\mu_0} \left( \mathbf{B} \mathbf{B} - \frac{B^2}{2} \mathbf{I} \right) \right] &= 0 \\ \frac{\partial \mathbf{B}}{\partial t} + \nabla \times (-\mathbf{u}_e \times \mathbf{B}) &= 0 \\ \frac{\partial e}{\partial t} + \nabla \cdot \left[ \mathbf{u} \left( e + p_{\perp} + \frac{\mathbf{B}^2}{2\mu_0} \right) + \mathbf{u} \cdot \left( (p_{\parallel} - p_{\perp}) \mathbf{b} \mathbf{b} - \frac{\mathbf{B} \mathbf{B}}{2\mu_0} \right) \right] &= 0 \end{aligned}$$

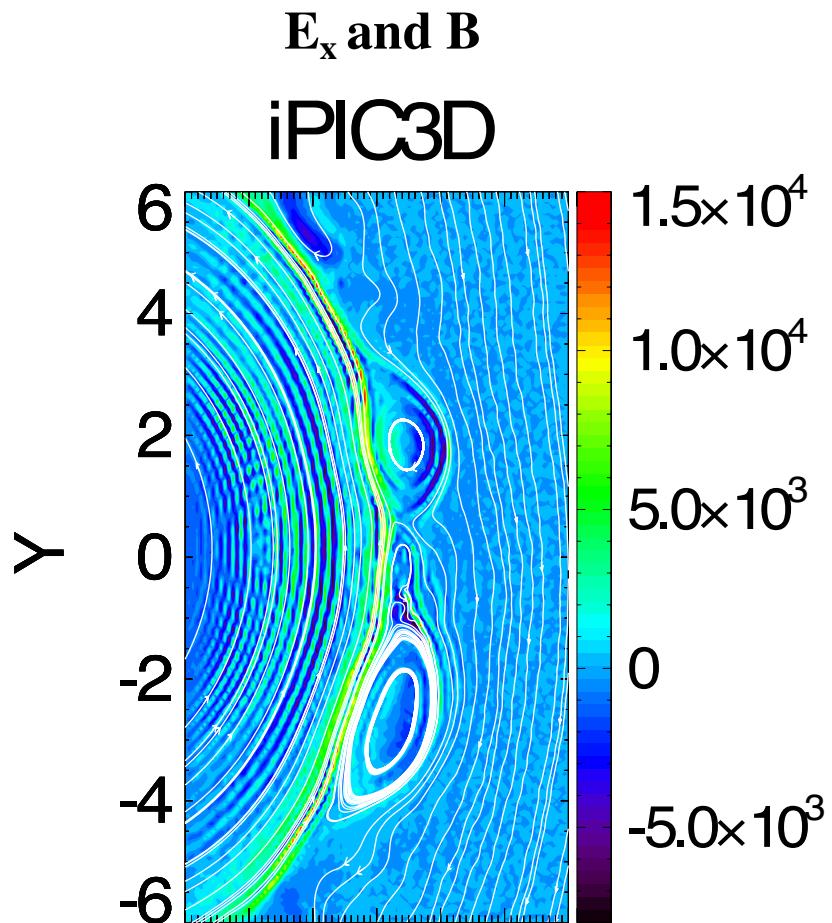
Particle-in-Cell:

$$\begin{aligned} \frac{\partial \mathbf{B}}{\partial t} &= -c \nabla \times \mathbf{E} \\ \frac{\partial \mathbf{E}}{\partial t} &= c \nabla \times \mathbf{B} - 4\pi \mathbf{J} \\ \frac{d\mathbf{p}}{dt} &= q(\mathbf{E} + \frac{\mathbf{V} \times \mathbf{B}}{c}) \end{aligned}$$

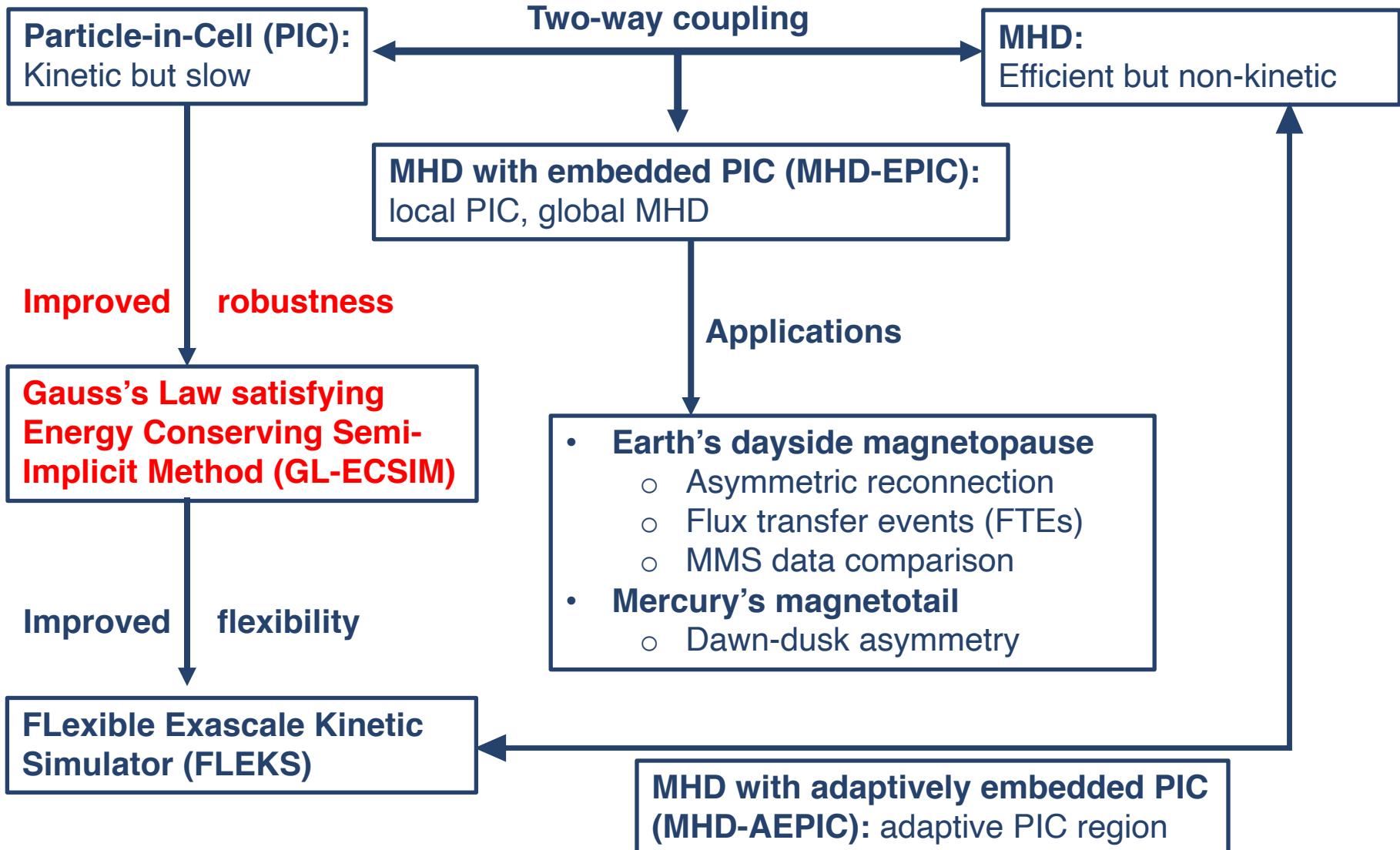
## ➤ Initialize PIC from MHD

- Calculate Electric field from Ohm's law:  $\mathbf{E} = -\mathbf{u} \times \mathbf{B} + \frac{\mathbf{J} \times \mathbf{B}}{q_e n_e}$
- Assume charge neutral:  $n_i = n_e$
- Ion and electron velocities can be obtained from the fluid momentum and current  $n(m_i \mathbf{u}_i + m_e \mathbf{u}_e) = \rho \mathbf{u}_{\text{MHD}}$        $n(q_i \mathbf{u}_i + q_e \mathbf{u}_e) = \mathbf{J}_{\text{MHD}}$
- Pressure  
Solve the electron pressure equation:  $\frac{\partial p_e}{\partial t} + \nabla \cdot (p_e \mathbf{u}_e) = (\gamma - 1)(-p_e \nabla \cdot \mathbf{u}_e)$
- Generate macro-particles from Maxwellian distribution

# Numerical Oscillations in iPIC3D



# Outline



# General Idea of ECSIM (Lapenta 2017)

Maxwell equations are solved implicitly on the grid:

$$\nabla \times \mathbf{E}^{n+\theta} + \frac{1}{c} \frac{\mathbf{B}^{n+1} - \mathbf{B}^n}{\Delta t} = 0 \quad \mathbf{E}^{n+\theta} = (1 - \theta)\mathbf{E}^n + \theta\mathbf{E}^{n+1}$$
$$\nabla \times \mathbf{B}^{n+\theta} - \frac{1}{c} \frac{\mathbf{E}^{n+1} - \mathbf{E}^n}{\Delta t} = \frac{4\pi}{c} \bar{\mathbf{J}} \quad \mathbf{B}^{n+\theta} = (1 - \theta)\mathbf{B}^n + \theta\mathbf{B}^{n+1}$$

The unknowns are at time level  $n+\theta$ . The current  $\bar{\mathbf{J}}$  at time level  $n+1/2$  is also calculated implicitly using  $\mathbf{E}^{n+\theta}$  and a mass matrix.

The particle velocity update uses  $\bar{\mathbf{J}}$ .

**For  $\theta=1/2$  the total energy (particles+fields) is exactly conserved**

This means that instabilities requiring energy cannot grow.

**Improving efficiency** (borrowing idea from iPIC3D)

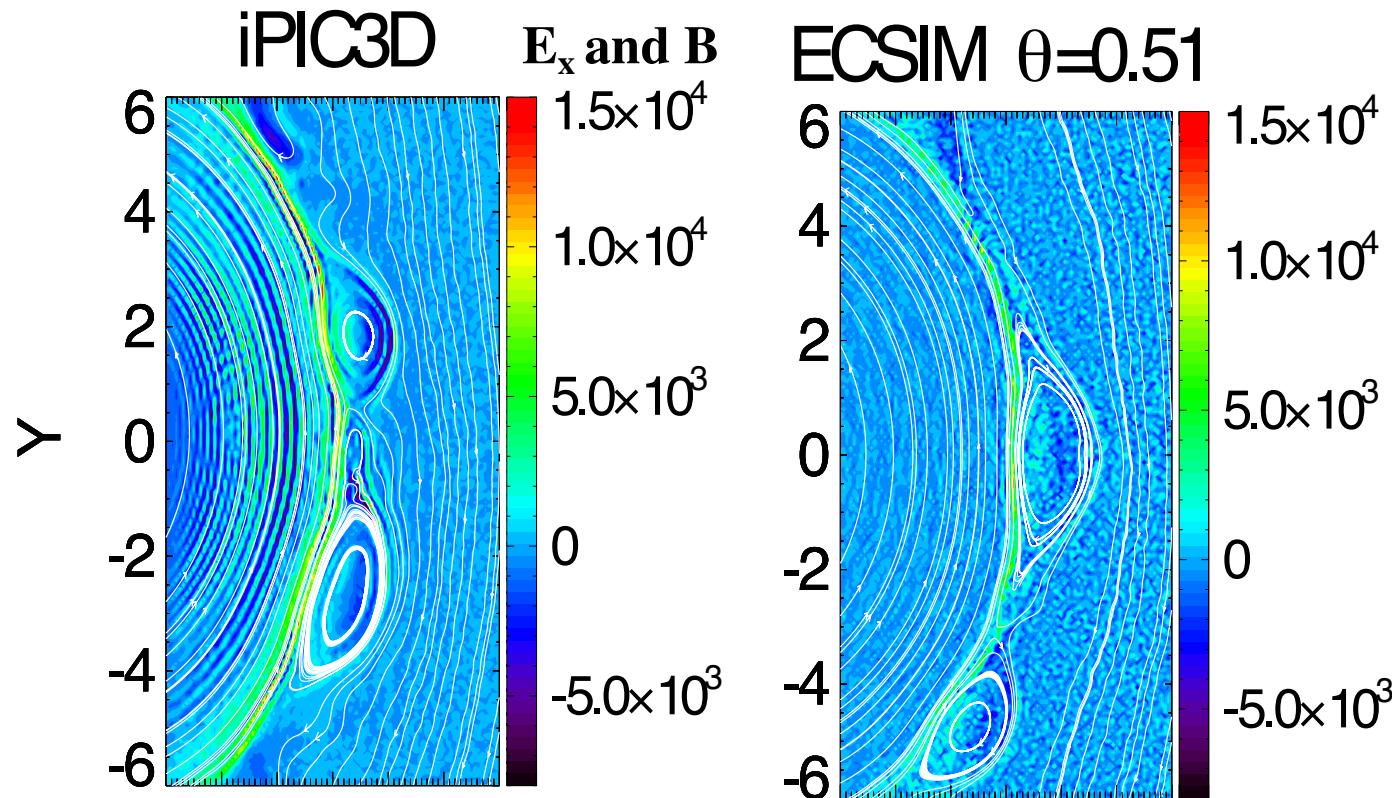
The above system can be combined into a numerically identical equation for  $\mathbf{E}^{n+\theta}$ :

$$\mathbf{E}^{n+\theta} + \delta^2 [\nabla(\nabla \cdot \mathbf{E}^{n+\theta}) - \nabla^2 \mathbf{E}^{n+\theta}] = \mathbf{E}^n + \delta(\nabla \times \mathbf{B}^n - \frac{4\pi}{c} \bar{\mathbf{J}})$$

where  $\delta=c\theta\Delta t$ . This is faster to solve than the original system of equations.

$\mathbf{B}^{n+1}$  is obtained from the induction equation explicitly.

# Comparison between iPIC3D and ECSIM



## ➤ 2D magnetopause simulation

- ECSIM with  $\theta=0.51$  produces good results. Small energy dissipation.

# Gauss's Law in ECSIM

- There is no mechanism to make the solution obey Gauss's law  $\rho_c = \nabla \cdot E / 4\pi$

## ➤ Remedies

- iPIC3D uses Marder's scheme to diffuse away the error, but using similar approach in ECSIM destroys energy conservation and improvements compared to iPIC3D.
- Poisson solver to project the electric field

$$\nabla^2 \phi = \nabla \cdot E^{n+1} - 4\pi \rho_c^{n+1}$$

$$E_{new}^{n+1} = E^{n+1} - \nabla \phi$$

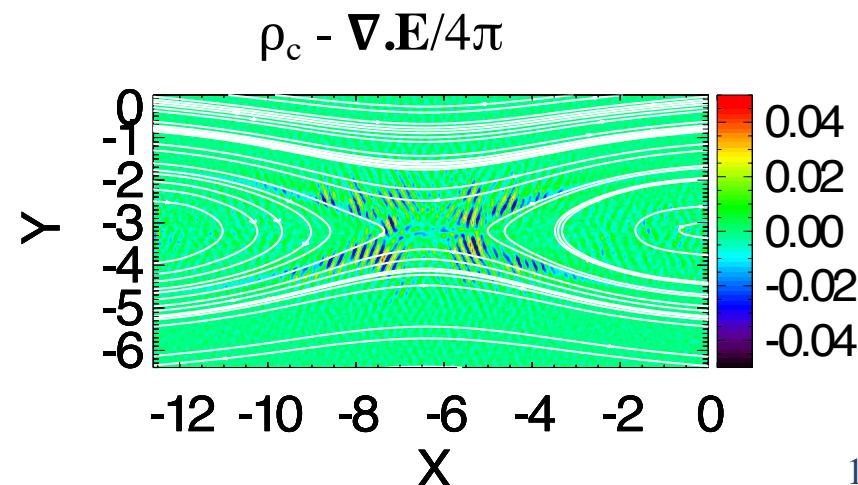
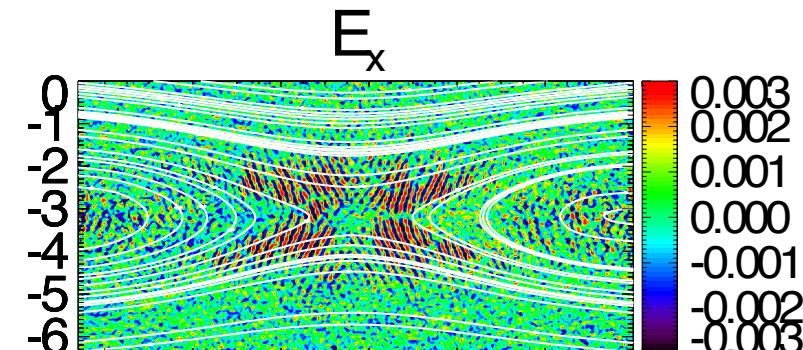
eliminates the error in Gauss's law, but energy is not conserved and results look similar to iPIC3D.

## ➤ What to do?

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi \mathbf{J}$$

ECSIM  $\theta=0.51$



# Move Particles to Satisfy Gauss's Law

## ➤ Why?

- Energy is still conserved: fields and velocity are unchanged, only particle positions are modified.
- Integration of particle trajectories is approximate, so why not?

## ➤ How?

- After completing the ECSIM update **adjust particle positions** to satisfy Gauss's law while **minimizing the total displacement: Lagrange multiplier method**. (Chen and Toth, 2019, JCP)

## ➤ Our goal is to make the error zero at time level n+1

- Time staggering (fields are at integer levels, particles at half levels) we need to interpolate:

$$\rho_c^{n+\frac{1}{2}} = \sum_p q_p W(\mathbf{x}_p^{n+\frac{1}{2}} - \mathbf{x}_c)$$

$$\rho_c^{n+1} = \gamma \sum_p q_p W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} + \Delta \mathbf{x}_p - \mathbf{x}_c) + (1 - \gamma) \rho_c^{n+\frac{1}{2}} = \frac{1}{4\pi} \nabla \cdot \mathbf{E}^{n+1}$$

- $\gamma=0.5$  gives second order, but  $\gamma=0.51$  works better.

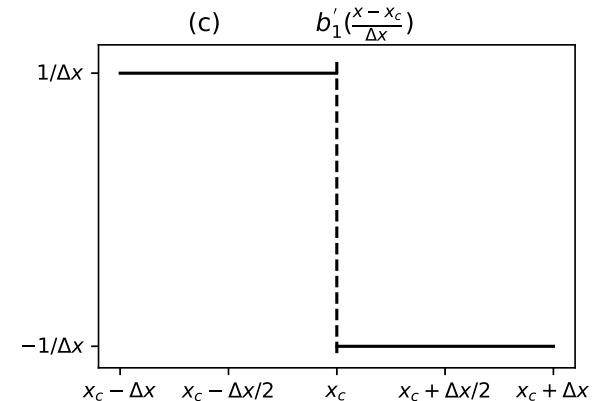
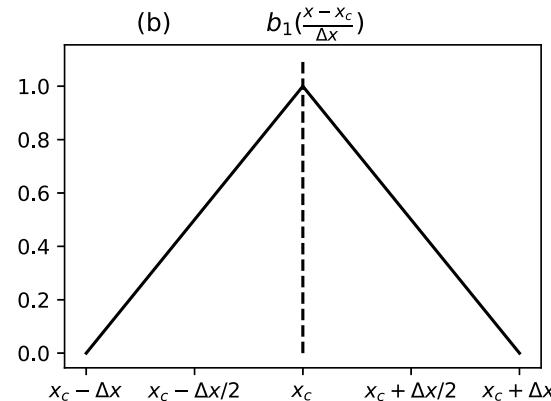
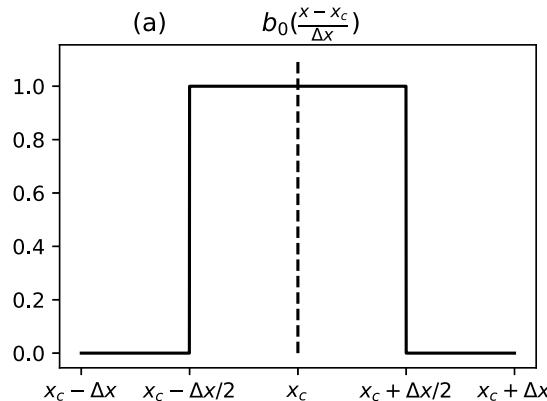
# Linearization

Linearization of charge density in  $\Delta\mathbf{x}_p$  is necessary for efficient solver:

$$\rho_c^{n+1} = \gamma \sum_p q_p W\left(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} + \Delta\mathbf{x}_p - \mathbf{x}_c\right) + (1 - \gamma)\rho_c^{n+\frac{1}{2}} = \frac{1}{4\pi}\nabla \cdot \mathbf{E}^{n+1}$$

$$W\left(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} + \Delta\mathbf{x}_p - \mathbf{x}_c\right) = W\left(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c\right) + \nabla W\left(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c\right) \cdot \Delta\mathbf{x}_p + O((\Delta x)^2)$$

ECSIM uses  $b_1$  spline weight functions, but there is a singularity and non-linearity in 2D and 3D, so the linearization is approximate



# Solving with Lagrange Multipliers

Minimize  $f(\Delta \mathbf{x}_p) = \sum_p \frac{1}{2} (\Delta \mathbf{x}_p)^2 |q_p|^\alpha$  (we will use  $\alpha=1$ ) with the constraints

$$g_c(\Delta \mathbf{x}_p) \triangleq \sum_p q_p \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c) \cdot \Delta \mathbf{x}_p - S_c = 0$$

Construct Lagrange function with Lagrange multipliers  $\lambda_c$ :

$$L(\Delta \mathbf{x}_p, \lambda_c) = f(\Delta \mathbf{x}_p) - \sum_c \lambda_c g_c(\Delta \mathbf{x}_p)$$

Take derivatives of L with respect to  $\Delta \mathbf{x}_p$ :

$$\frac{\partial L}{\partial \Delta \mathbf{x}_p} = \Delta \mathbf{x}_p |q_p|^\alpha - \sum_c \lambda_c q_p \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c) = 0$$

Thanks to the linearization of  $g_c$ , this can be solved for  $\Delta \mathbf{x}_p$ . Substitute into  $g_c=0$ :

$$\frac{\partial L}{\partial \lambda_c} = \sum_p q_p \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c) \cdot \left[ |q_p|^{-\alpha} \sum_{c'} \lambda_{c'} q_p \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_{c'}) \right] - S_c = 0$$

Swap the sums for p and c' to obtain a linear system for  $\lambda_c$  with a "mass matrix":

$$\sum_{c'} \lambda_{c'} M_{cc'} = S_c \quad M_{cc'} = \sum_p |q_p|^{2-\alpha} \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c) \cdot \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_{c'})$$

# Non-linear and Alternative Schemes

## ➤ Iterate over the linear solver (~3 times):

- Solve the linear system

$$\sum_{c'} \lambda_{c'} M_{cc'} = S_c$$

- for  $\lambda_c$ , then calculate

$$\Delta \mathbf{x}_p = \sum_c \lambda_c |q_p|^{1-\alpha} \frac{q_p}{|q_p|} \nabla W(\tilde{\mathbf{x}}_p^{n+\frac{3}{2}} - \mathbf{x}_c)$$

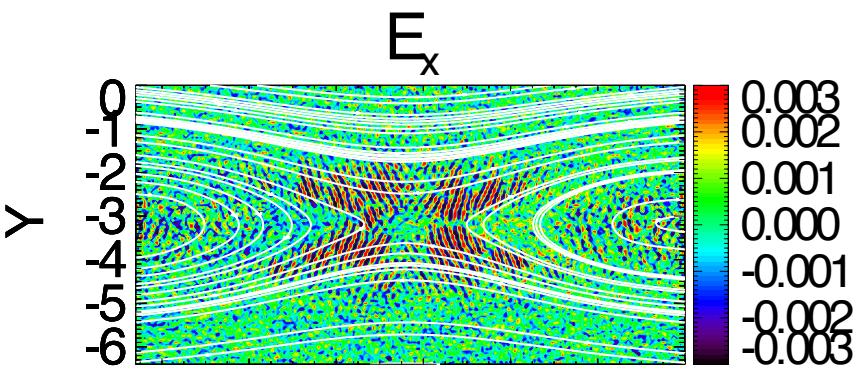
- Update particle positions, recalculate  $M_{cc'}$  and  $S_c$  and repeat

## ➤ Alternatives

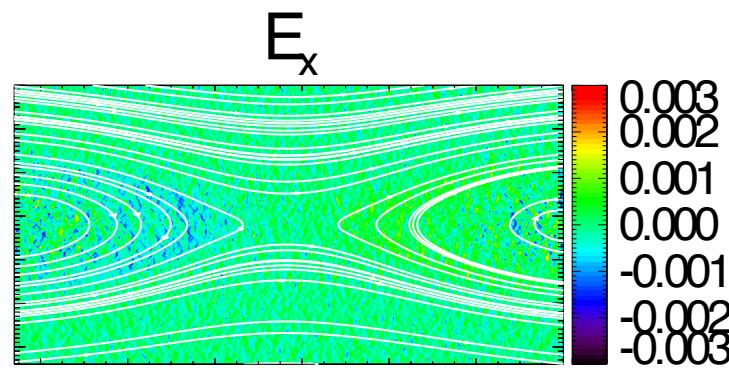
- Approximate solver that does not use a mass matrix
- Local “push” of particles to reduce error in Gauss’s law

# GEM Reconnection

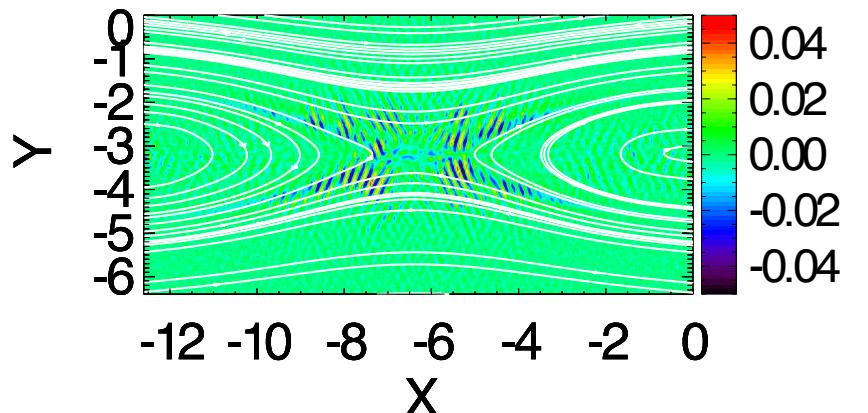
ECSIM  $\theta=0.51$



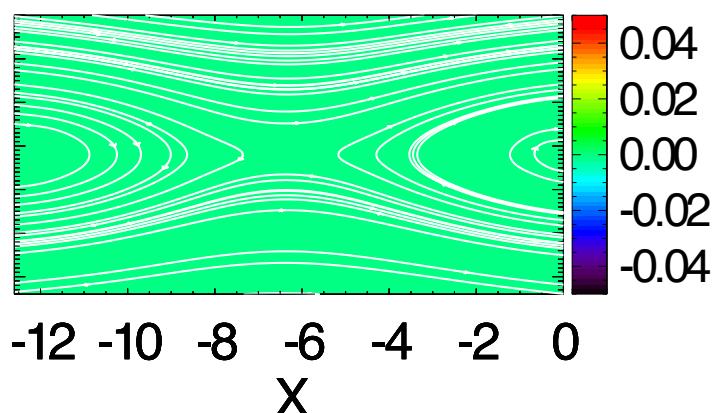
GL-ECSIM



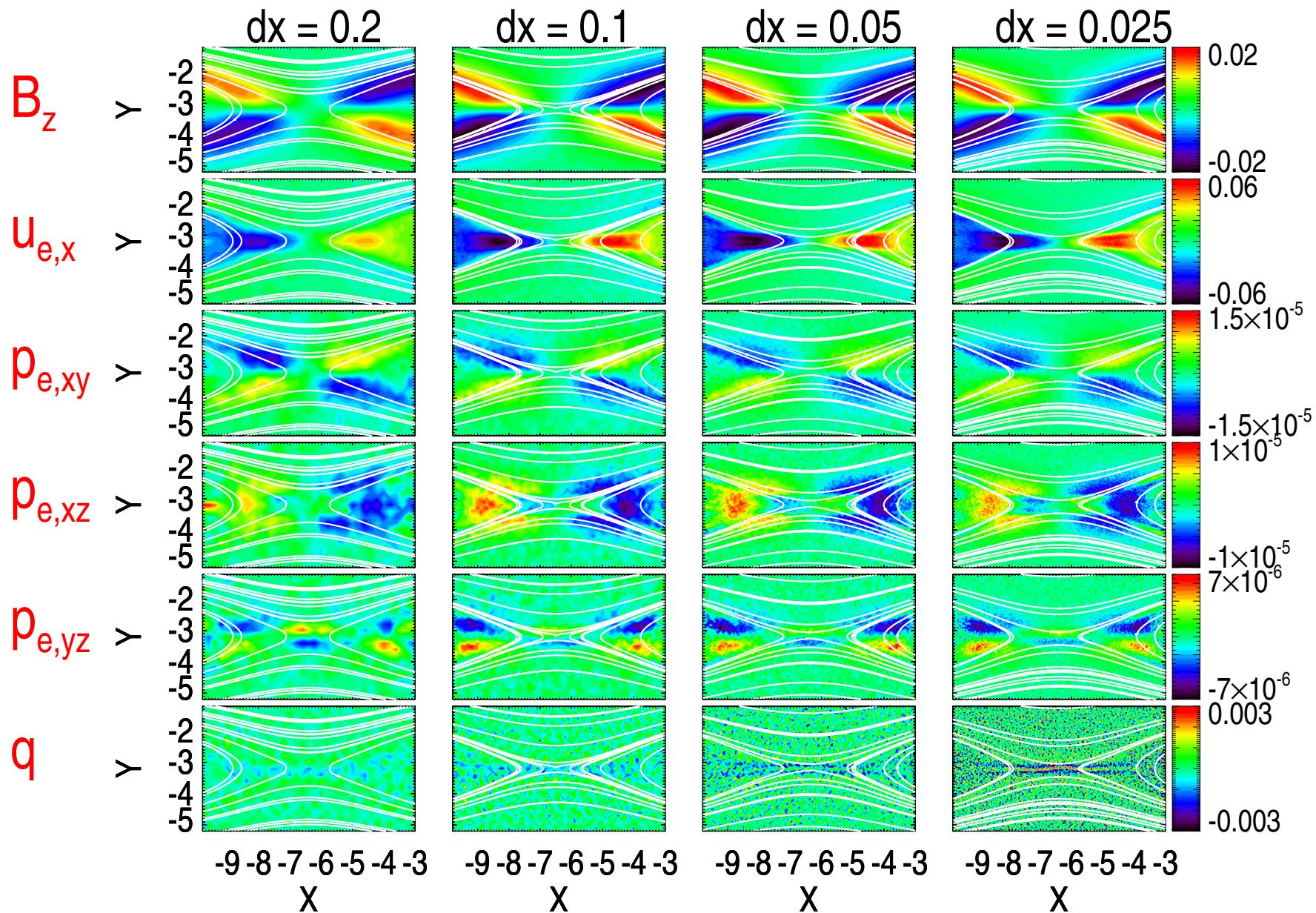
$\rho_c - \nabla \cdot E / 4\pi$



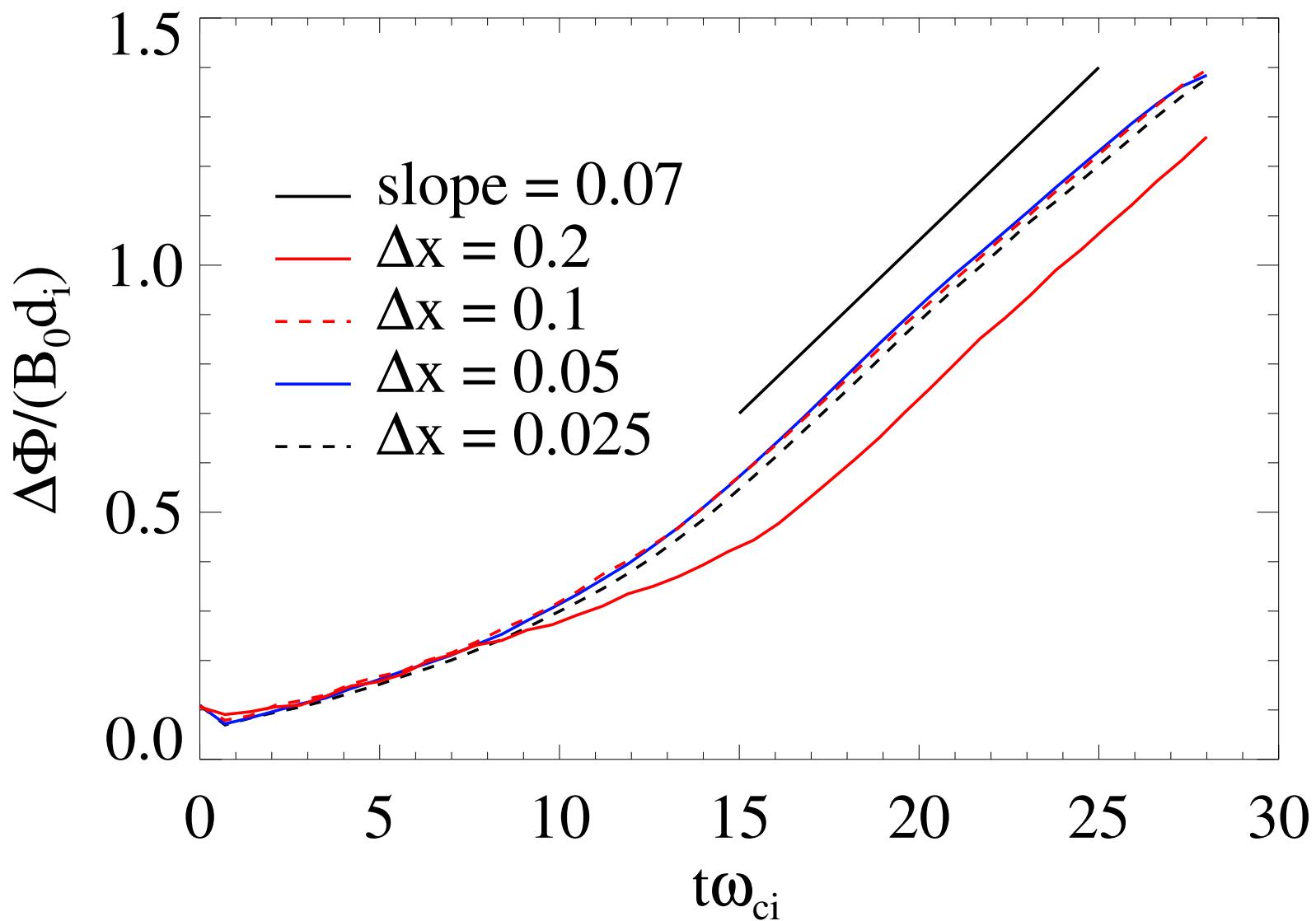
$\rho_c - \nabla \cdot E / 4\pi$



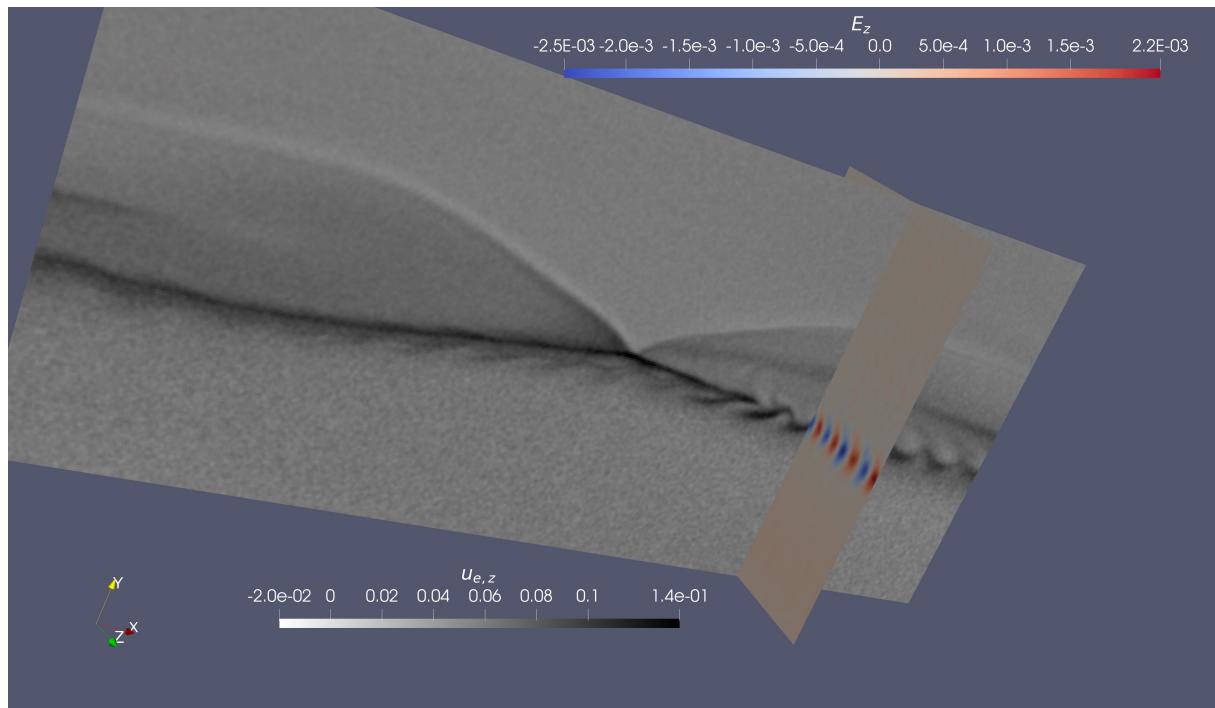
# Grid Convergence ( $d_i=1$ , $d_e=0.1$ )



## Grid Convergence ( $d_i=1$ , $d_e=0.1$ )



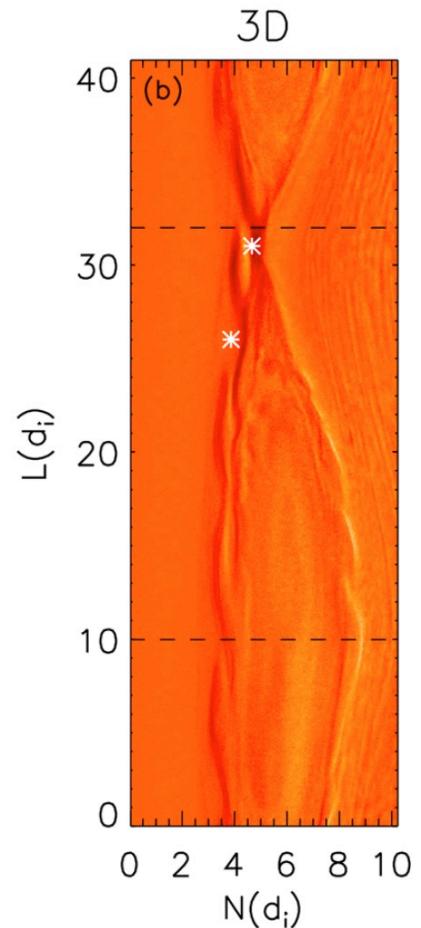
# GL-ECSIM vs. Explicit PIC



GL-ECSIM

$\Delta x = 1/16 d_i \approx 0.06 d_i$

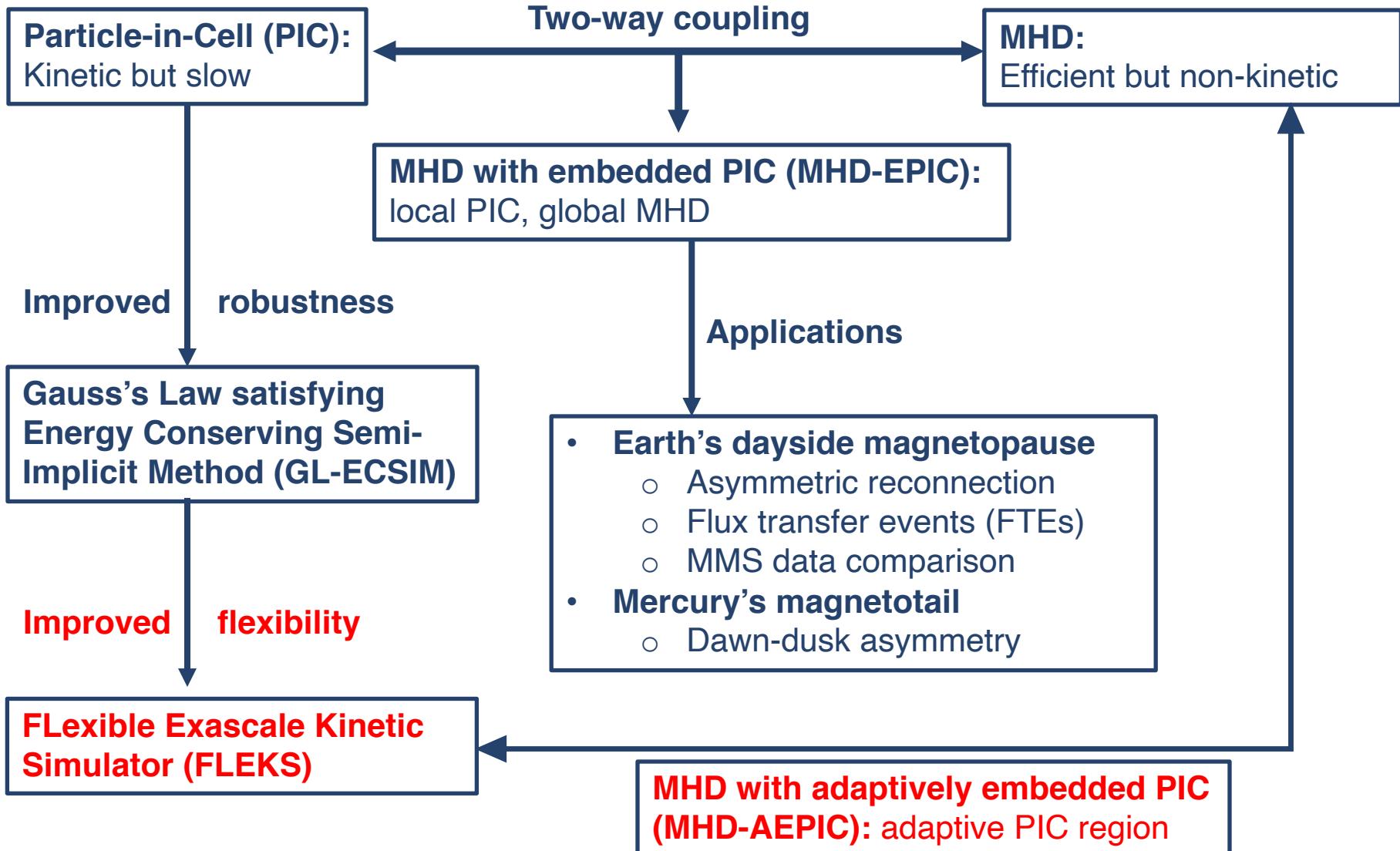
Electron current density



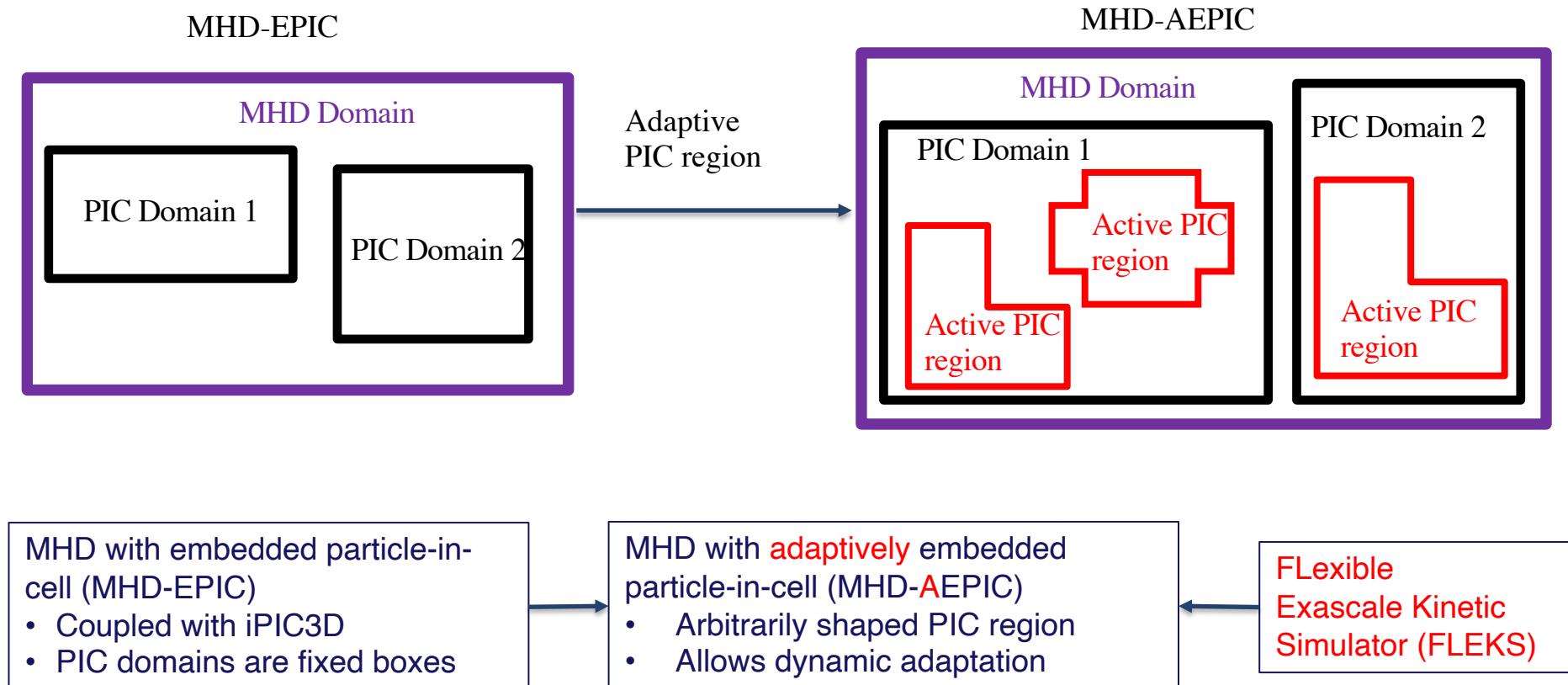
Explicit PIC (p3d).  
 $\Delta x \approx 0.02 d_i$

(Price et al., 2016, GRL)

# Outline

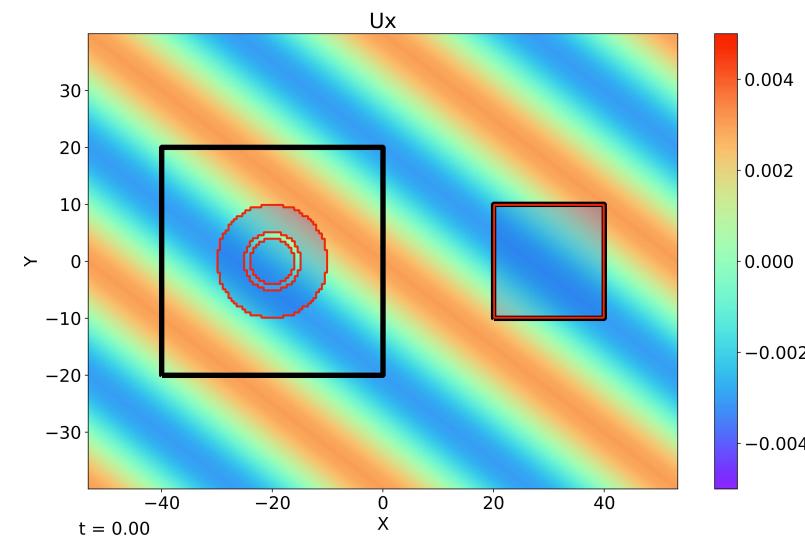


# From MHD-EPIC to MHD-AEPIC



# FFlexible Exascale Kinetic Simulator (FLEKS)

- FLEKS is a brand new semi-implicit PIC code
- PIC algorithm
  - Gauss's law satisfying energy-conserving semi-implicit particle-in-cell method (GL-ECSIM)
  - Particle resampling: splitting and merging
- PIC grid
  - Multiple independent PIC domains
  - Cartesian grid
  - Adaptation: PIC cells can be switched on/off
- Implementation
  - A third-party library AMReX provides high-performance parallel data structure
  - Parallelized with MPI. No OpenMP/GPU so far

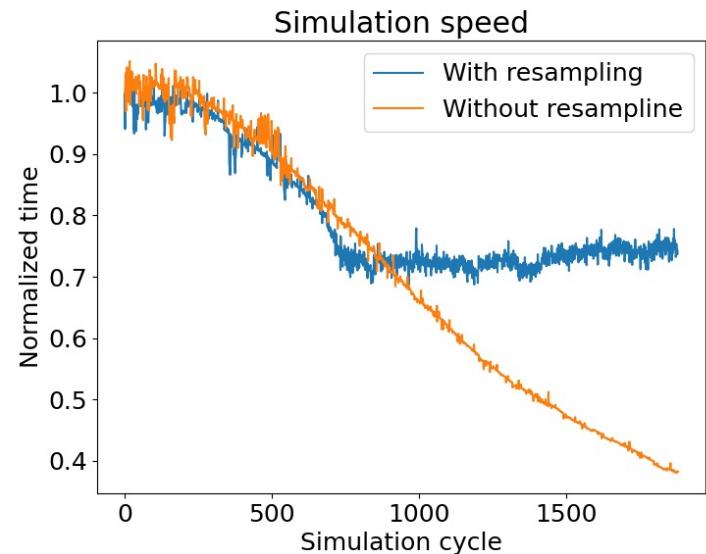


- PIC domains: black squares
- Active PIC regions: shadowed areas, enclosed by red lines

# Particle Resampling: Splitting and Merging

## ➤ Problem

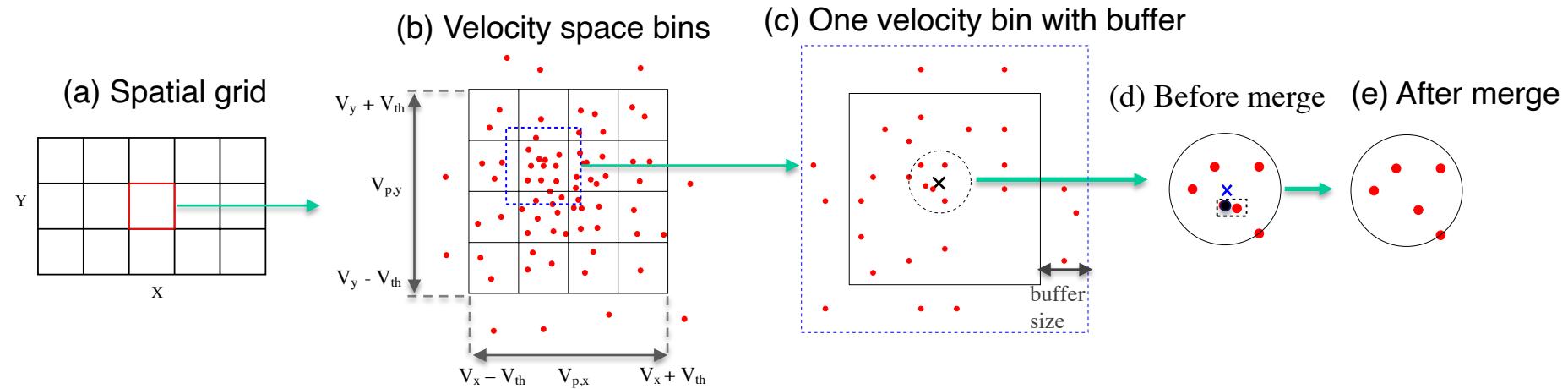
- Concentration of mass (macro-particles)
  - > load imbalance
  - > slows down the simulation
- Decreasing mass density (macro-particles)
  - > zero macro-particles in some cells
  - > reduces the simulation accuracy



## ➤ Solution

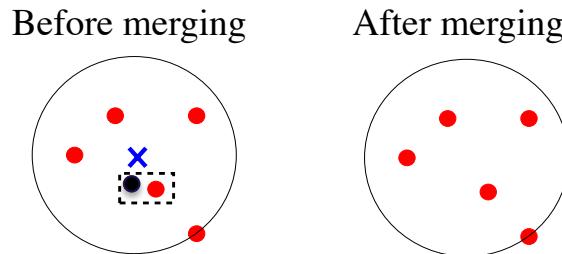
- Number of macro-particles per cell  $< N_{\text{low}}$  : split 1 particle into 2
- Number of macro-particles per cell  $> N_{\text{high}}$  : merge 6 particles into 5
  - Find 6 particles that are close to each other in the 6D phase space.
  - Merge 6 particles into 5. Conserve mass, momentum and energy.

# Particle Merging (1)



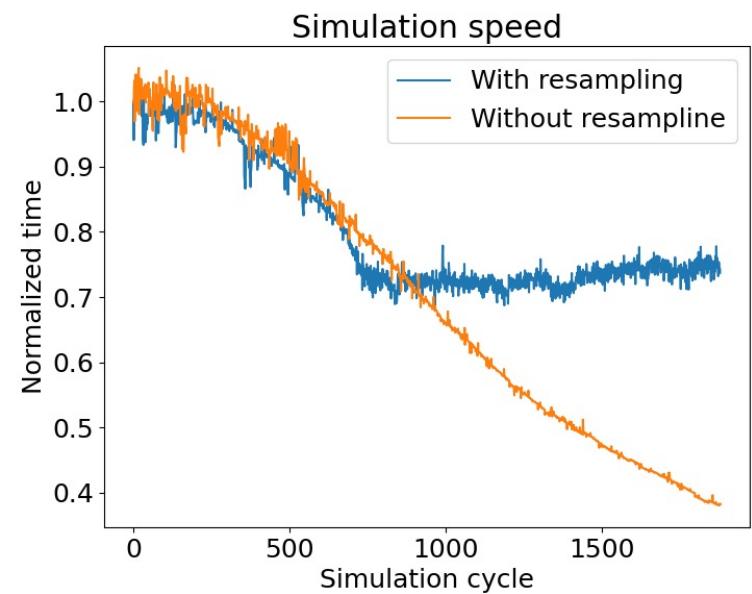
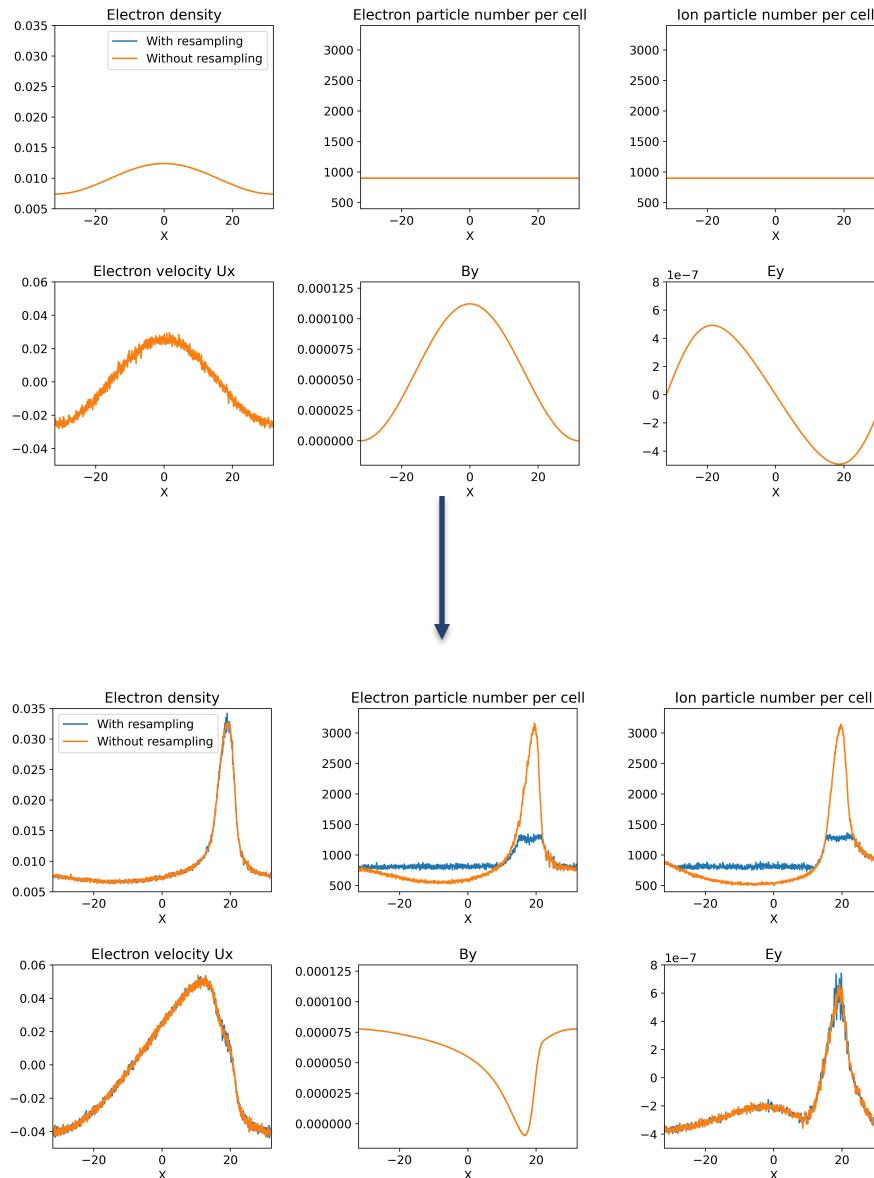
- **Step 1:** For each spatial cell (a) create a grid in the velocity space ranging from  $V-V_{th}$  to  $V+V_{th}$  and assign particles to velocity bins (b).
- **Step 2:** For each velocity bin, calculate the center of all the particles (black cross in panel c), and find the 6 particles closest to the center.
- **Step 3:** Find the 6D center (blue cross in panel d) of these 6 particles, the 6D distance of a particle to the center should NOT exceed a threshold value.
- **Step 4:** Merge 6 particles into 5

## Particle Merging (2)

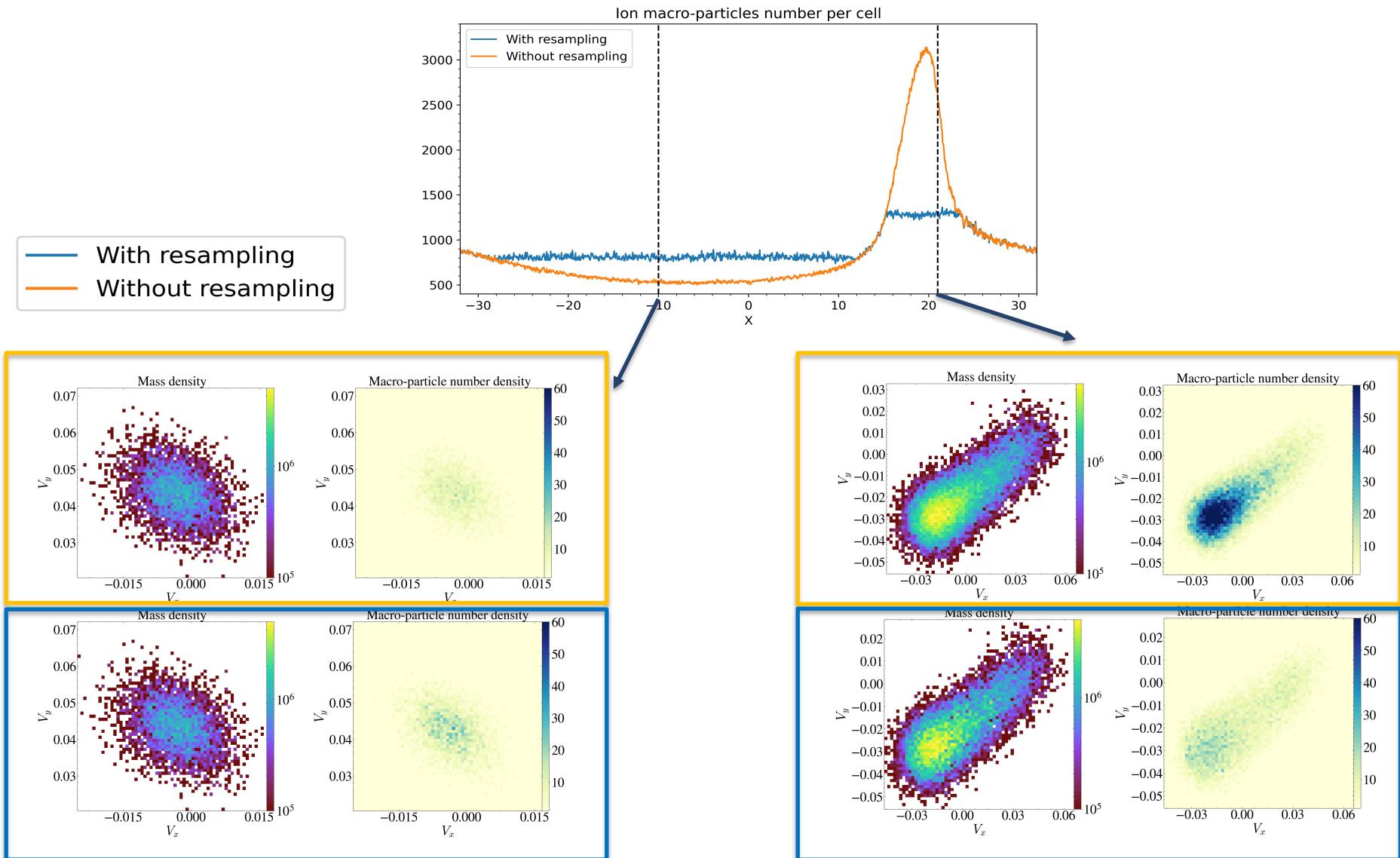


- From these 6 particles, find a pair that is closest to each other. Remove the lighter one and re-distribute its mass to the other 5 particles by changing their weights. These 5 particles keep their original positions and velocities. The total mass, momentum, and energy are conserved after merging by properly adjusting the weights:
  - 5 conserved quantities, needs 5 new particles.
  - A physical meaningful solution may not always exist.
  - This algorithm allows merging N particles into 5, where  $N \geq 6$
- Why not change velocities and merge into fewer particles?  
The distribution function could change drastically.

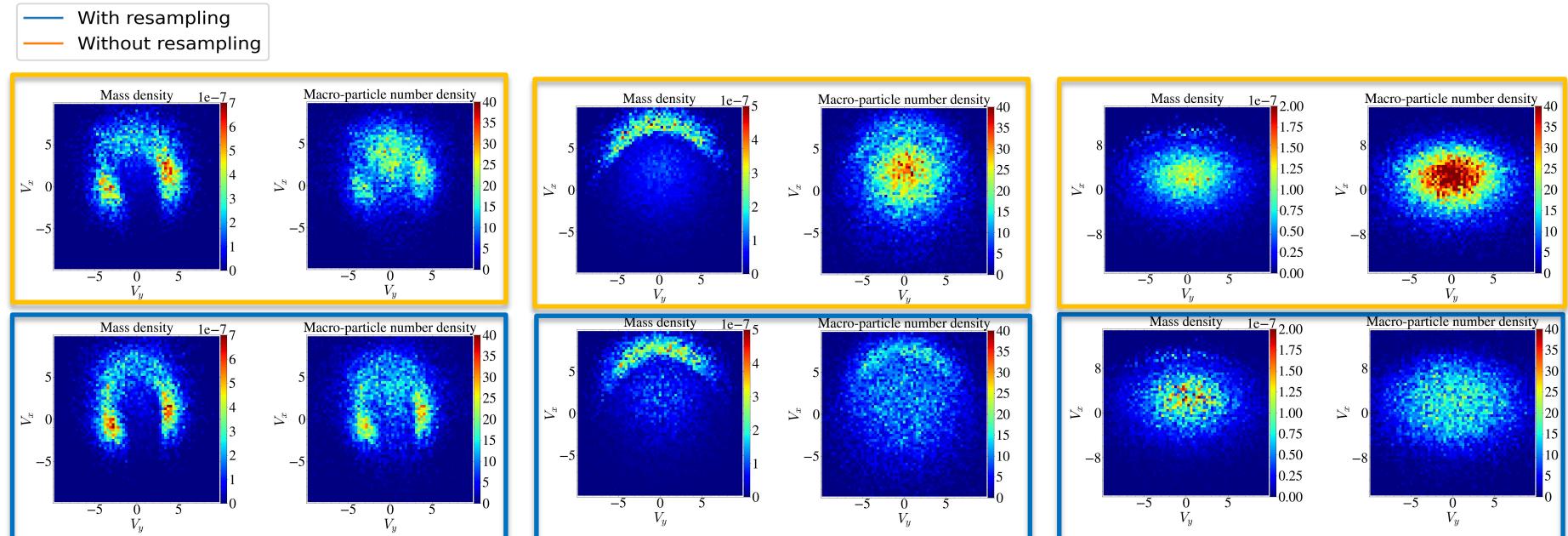
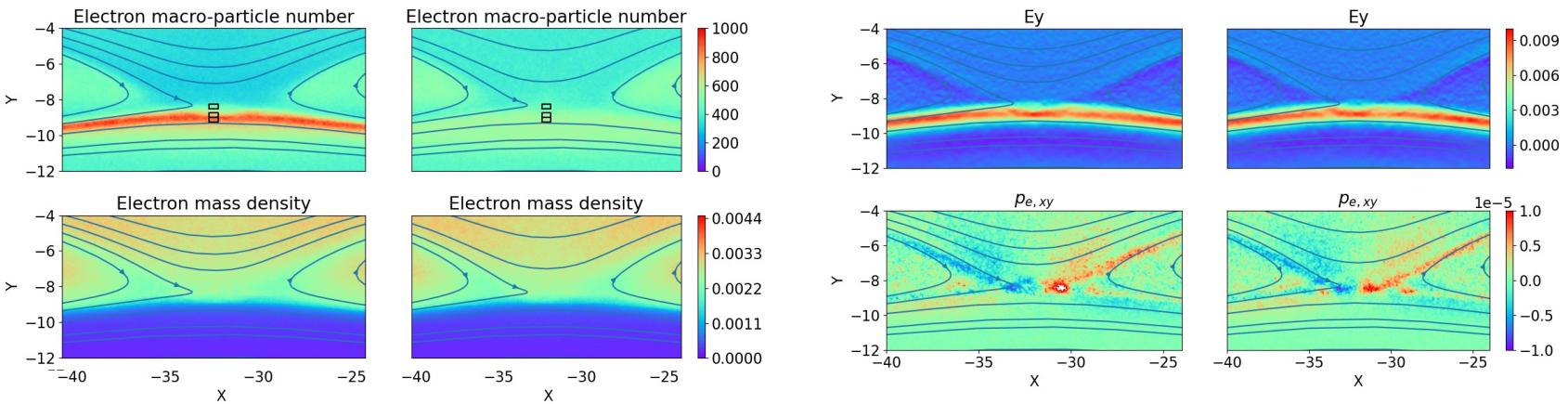
# Particle Resampling Test: 1D Magnetosonic Wave



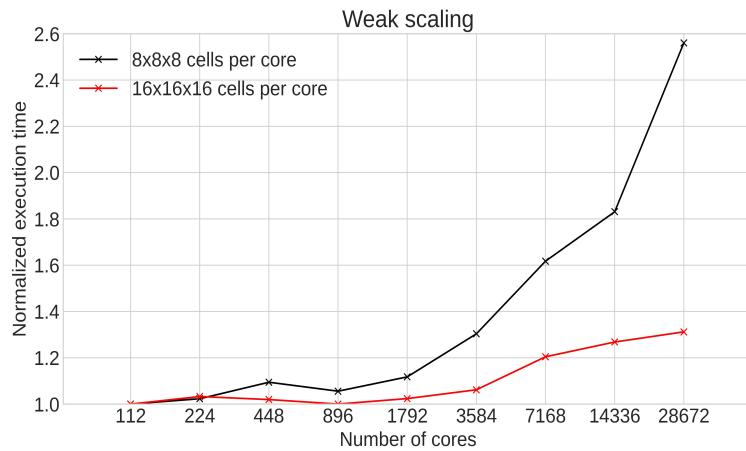
# Magnetosonic Wave Test Particle Distributions



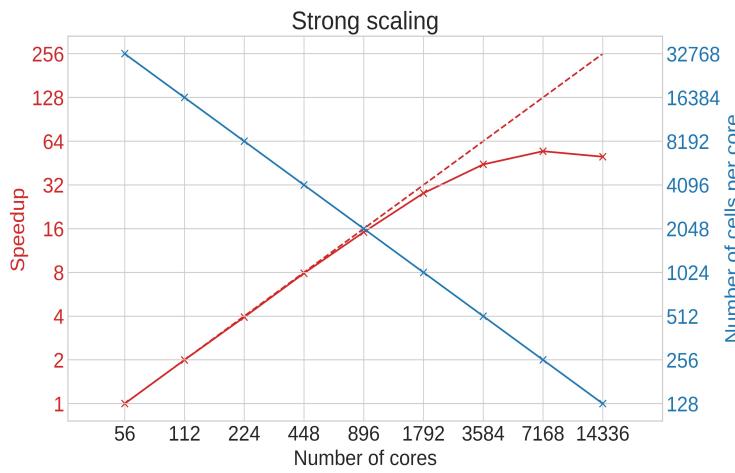
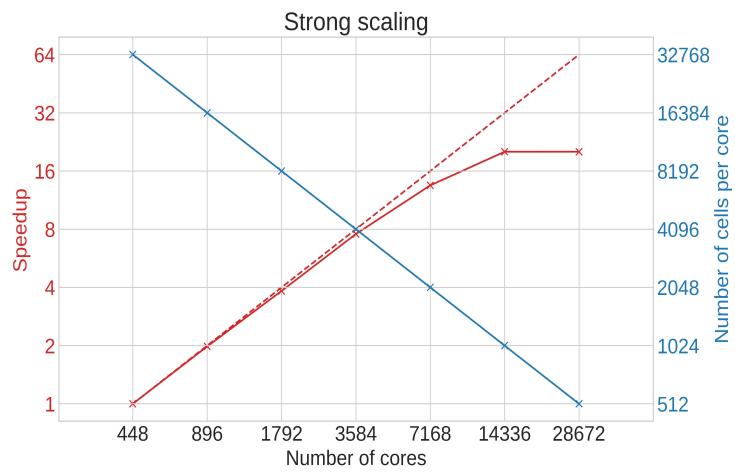
# Particle Resampling Test: 2D Asymmetric Magnetic Reconnection



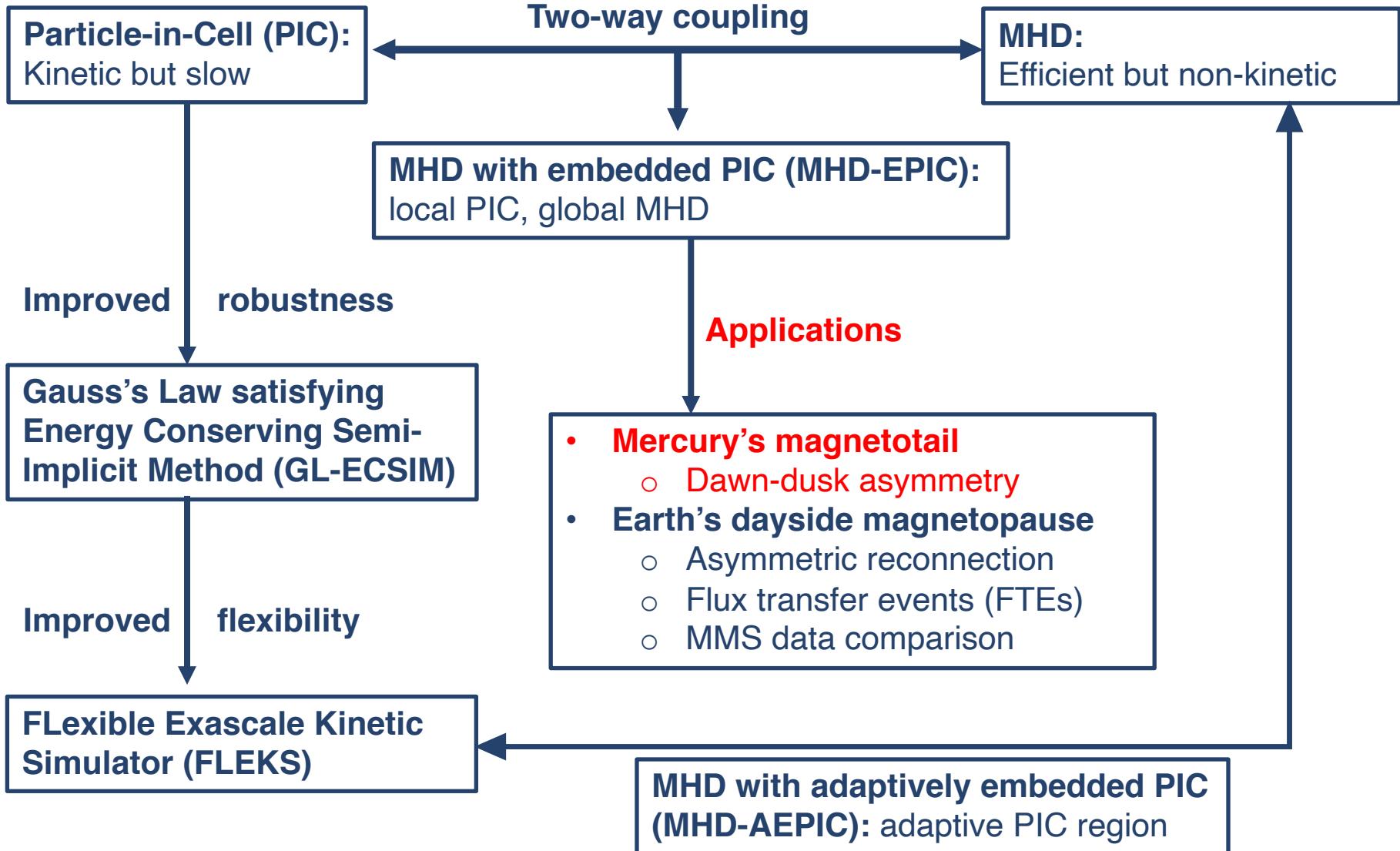
# Performance of FLEKS



- Run on Frontera, which ranks 8<sup>th</sup> on Top 500. 23.5 PetaFLOPS peak performance
- In general, FLEKS can efficiently use ~30k cores, which is about 7% of the whole machine
- Exascale? OpenMP or/and GPU are needed



# Outline



# MHD-EPIC Simulation of Mercury's Magnetotail

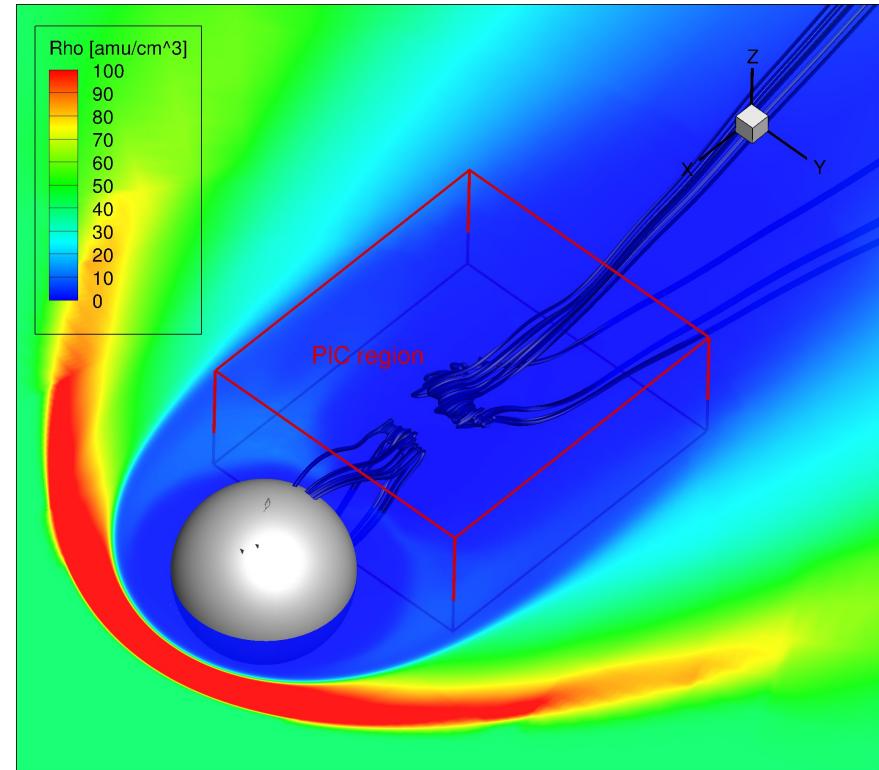
## ➤ Why Mercury?

- Mercury's magnetosphere is smaller than that of Earth
- Magnetotail ion scale can be resolved without scaling
- The kinetic effects may have global effects: **dawn-dusk asymmetries**

## ➤ Simulation

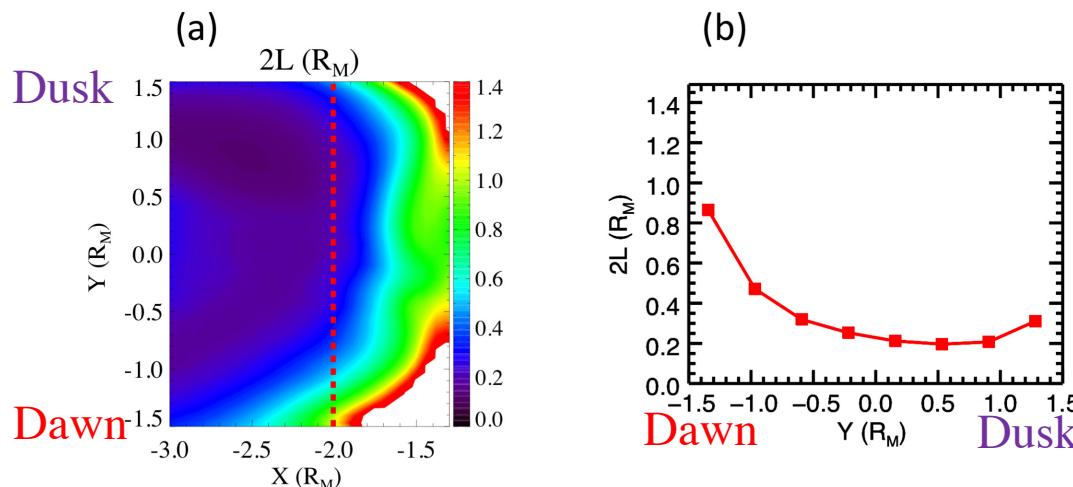
(Chen et al., 2019, JGR)

- PIC covers magnetotail
- $\Delta x = 1/32 R_m \sim 1/5 d_i$
- 300 s long  $\sim 2$  Dungey cycles
- $B_y = 0$  and  $u_y = 0$  in the solar wind
- IMF condition:  
 $B = [-17.4, 0, -8.5] \text{ nT}$



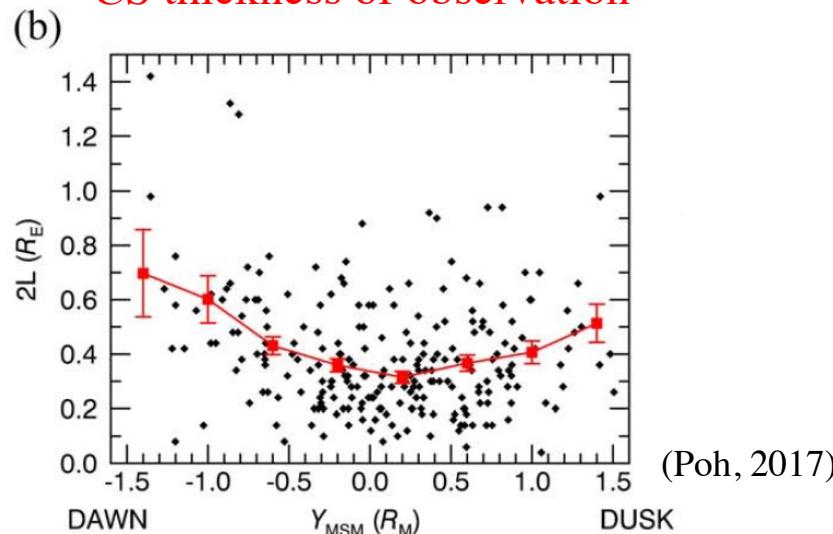
# Current-Sheet (CS) Thickness

CS thickness of MHD-EPIC simulations



- Cross-tail current-sheet is thicker on the **dawnside (-Y)** than the **duskside (+Y)**

CS thickness of observation



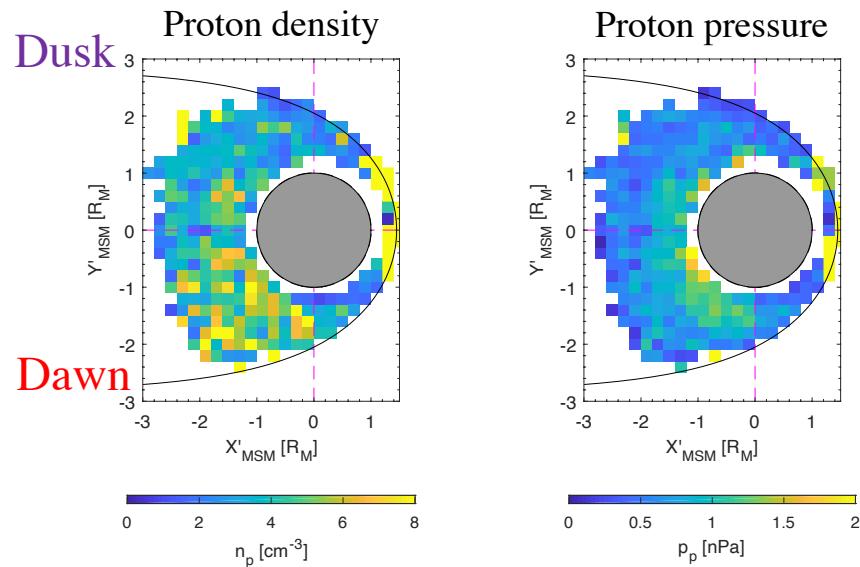
- Potential explanation: Hall effect transport magnetic flux and plasma from **duskside** to **dawnside (Lu, 2016)**

# Plasma Inside Current-Sheet

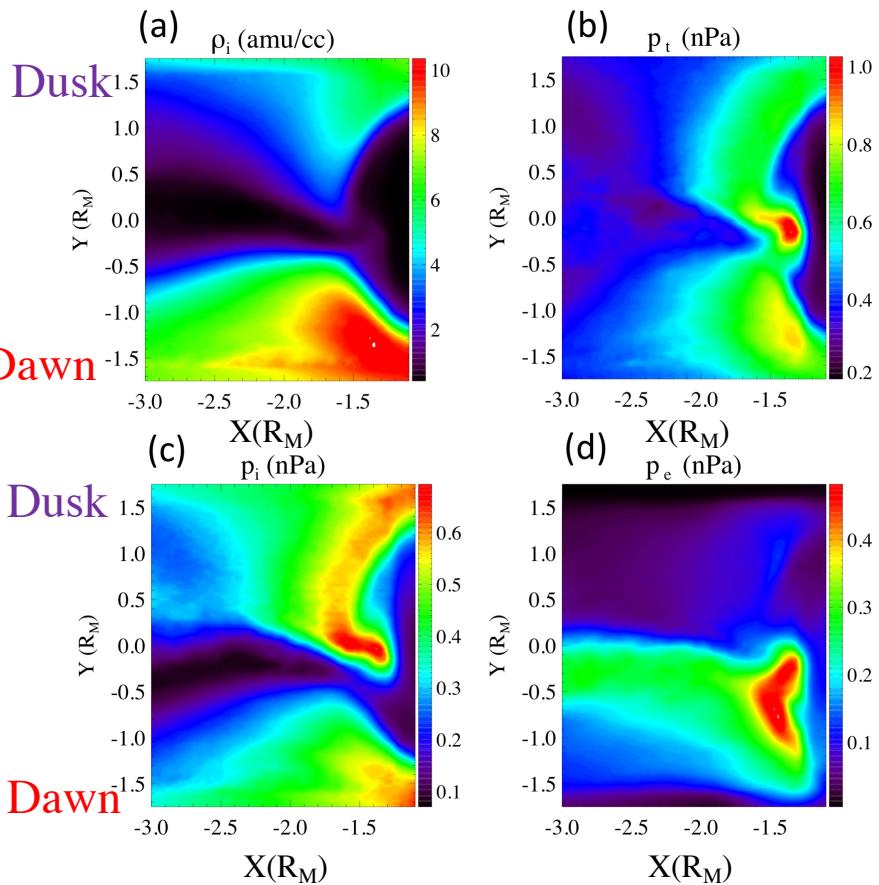
## ➤ Inside the current-sheet:

- Dawnside density is about twice of the duskside density
- Simulation ion pressure does not show significant dawn-dusk asymmetry
- Dawnside electron pressure is much higher than the duskside (no electron pressure observation)

MESSENGER observation

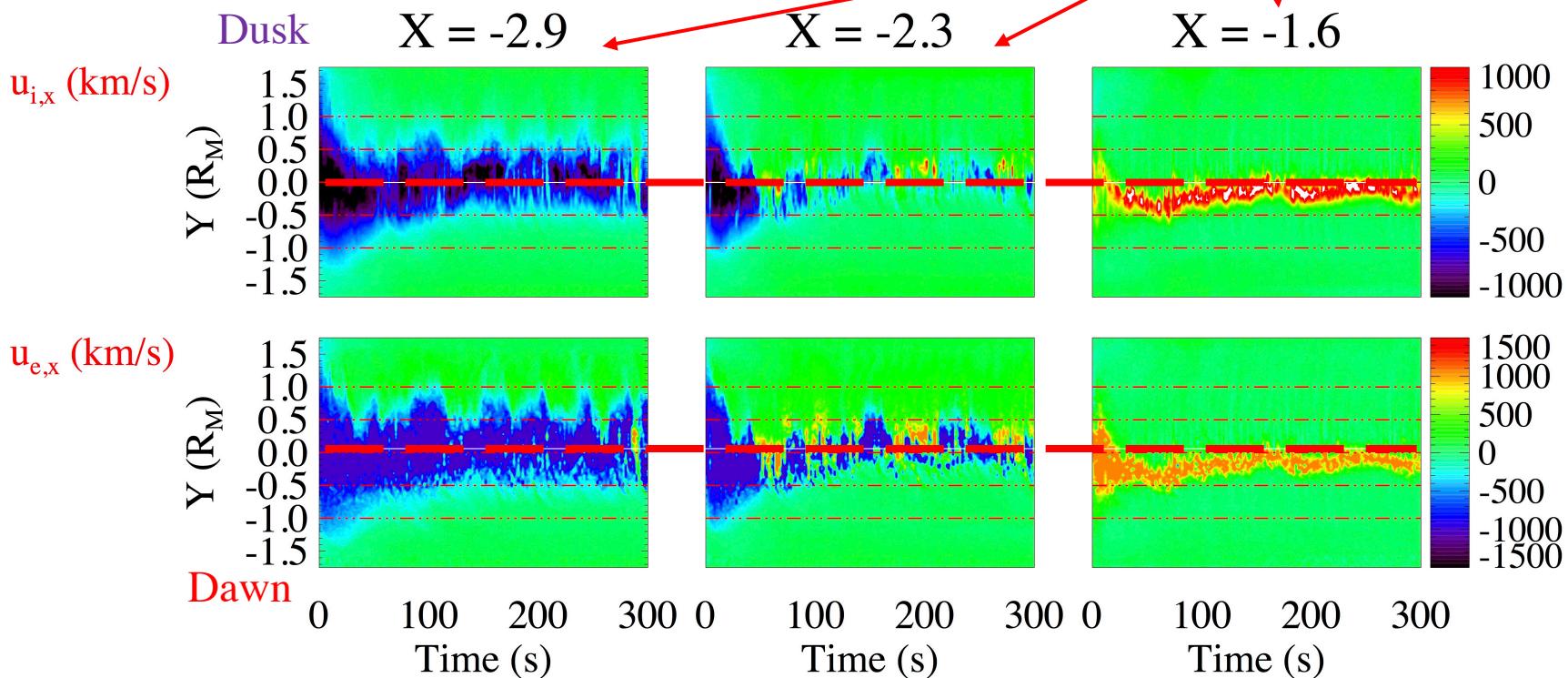


MHD-EPIC simulation



# Dawn-dusk Asymmetry of Reconnection Jets

- In this simulation, the reconnection jets do no show significant dawn-dusk asymmetry near the X-lines or on the tail-side
- The planetward jets prefer the **dawnside (-Y)**



# Dawnward Motion of Dipolarization Event

- Inside CS, electrons move **dawnward (-Y)**
- DF is also moving **dawnward (-Y)**

➤ **Inside the current-sheet, ions move duskward, but fast ion flows still prefer downside.**

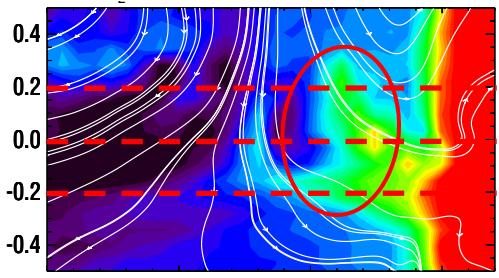
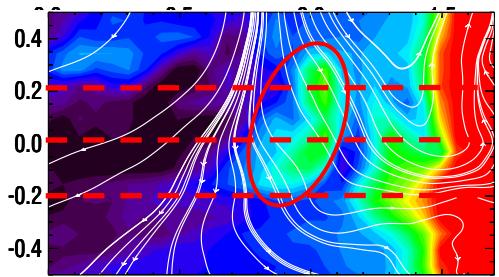
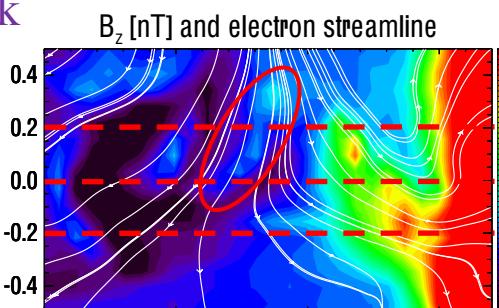
Maybe because fast ion flows are accelerated by dipolarization events

$t = 243.5 \text{ s}$

$t = 244.5 \text{ s}$

$t = 245 \text{ s}$

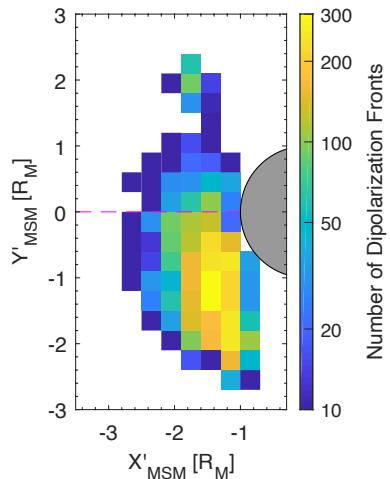
Dusk



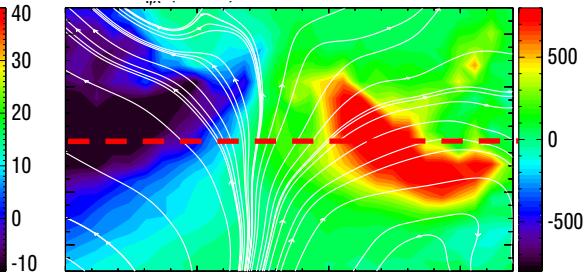
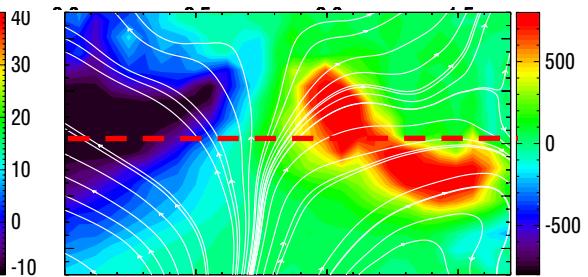
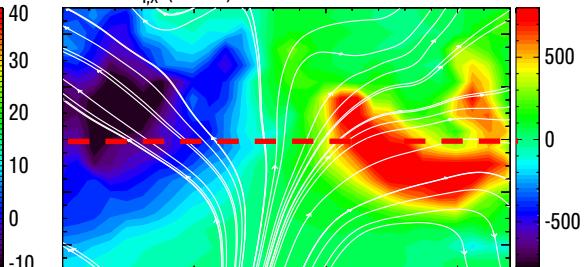
Dawn

-3.0 -2.5 -2.0 -1.5

X

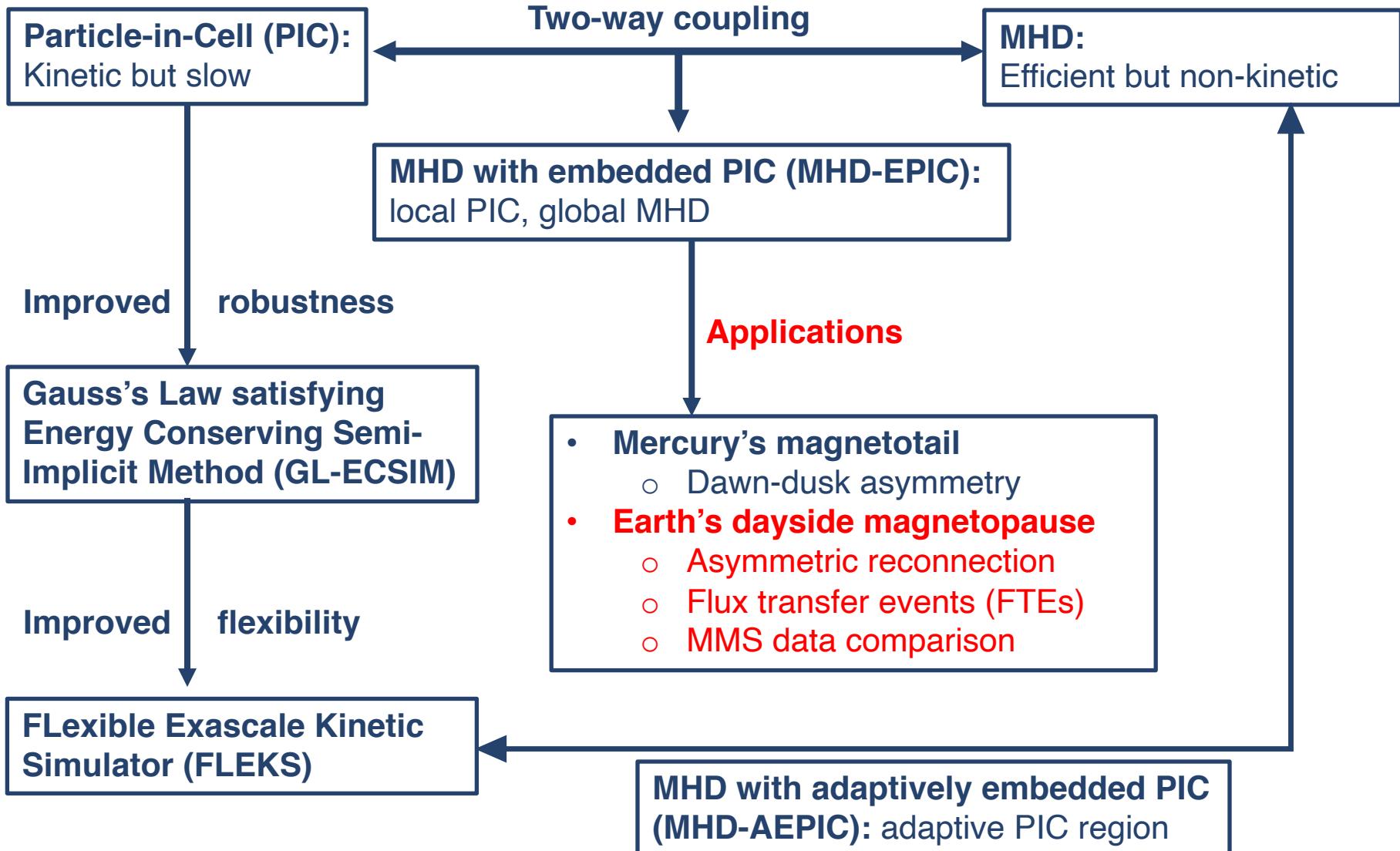


$u_{i,x}$  (km/s) and ion streamline



40

# Outline



# Introduction: Incorporate Kinetic Physics into Global Simulations

## ➤ Hybrid models

- Fluid massless electrons + kinetic (PIC or Vlasov solver) ions
- Local kinetic + global fluid -> MHD-EPIC

## ➤ Extend MHD

- Five-moment (isotropic pressure), six-moment (anisotropic pressures), and ten-moment (pressure tensors)

Earth's magnetosphere:

$\lambda_D \sim 100$  m;  $d_e \sim 5$  km;  $d_i \sim 100$  km;

Magnetosphere size  $\sim 10R_E \sim 10^5$  km

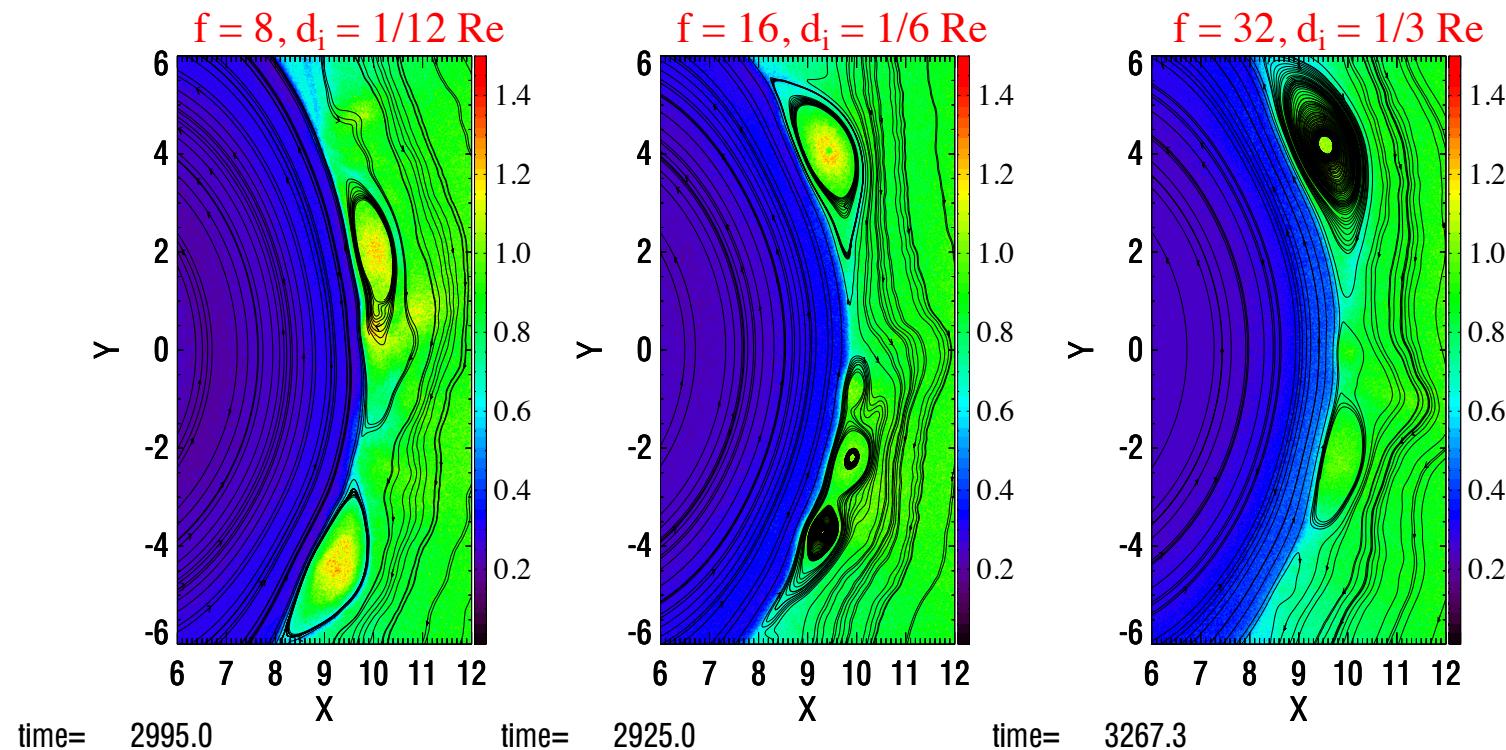
## ➤ Reduce the separation between scales

- Reduce speed of light: reduce separation between  $\lambda_D$  and  $d_e$
- Reduce the  $m_i/m_e$  ratio: reduce separation between  $d_e$  and  $d_i$
- **Reduce  $q_i/M_i$  to increase  $d_i$**

$$d_i = c \sqrt{\frac{M_i}{4\pi n q_i}}$$

# Changing the Kinetic Scales

- Reduce  $q_i/M_i$  by a scaling factor  $f$  to increase  $d_i$ , while the number density, pressure, thermal velocity do not change.  
(Toth et al., 2017, JGR)



# 3D MHD-EPIC Simulation Set Up

## ➤ Physical parameters

- Use  $f=16$  so  $d_i \sim 1/6 Re$ : doable
- Typical solar wind conditions:  
 $\rho = 5 \text{ amu/cm}^3$ ,  $U_x = -400 \text{ km/s}$ ,  $B = [0, 0, -5] \text{ nT}$

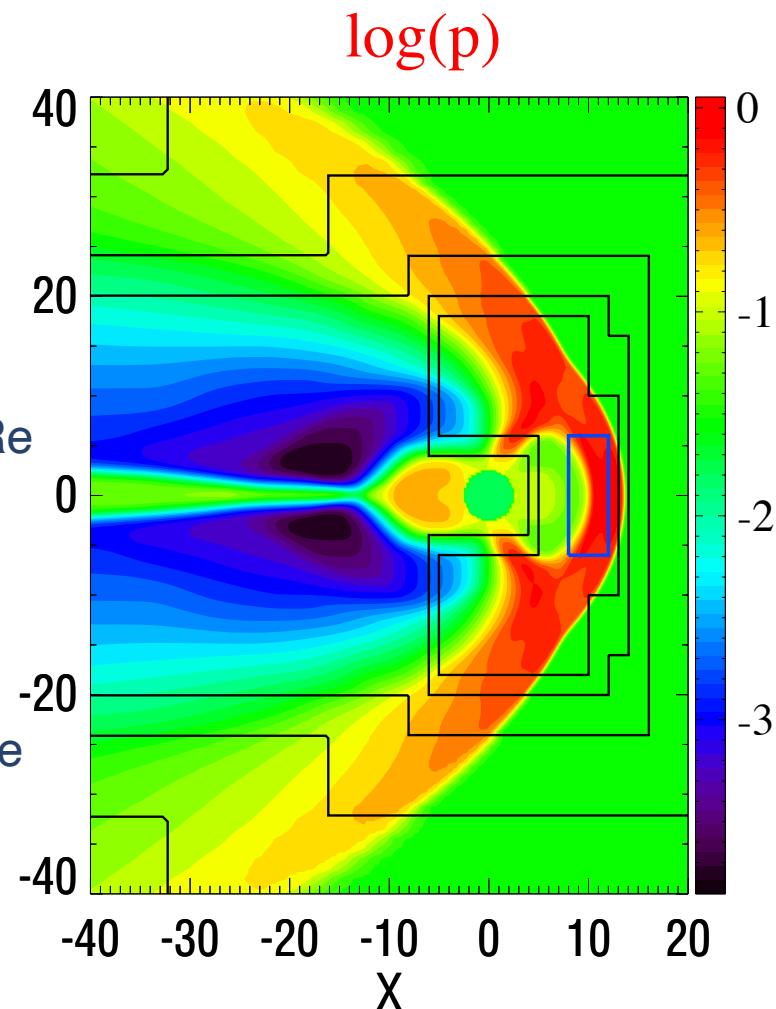
## ➤ Hall MHD with separate electron pressure equation

- MHD domain:  $-224 < x < 32$ ,  $-128 < y, z < 128Re$
- At the magnetopause  $\Delta x = 1/16Re$  ( $\sim 400 \text{ km}$ )  $\approx$
- MHD uses  $\sim 20\%$  CPU time

## ➤ PIC

- PIC domain:  $8 < x < 12$ ,  $-6 < y < 6$ ,  $-6 < z < 6Re$
- $\Delta x = 1/32Re$ : 5 cells per  $d_i$  (for  $f = 16$ )
- 216 particles per cell per species: 8B total
- Consuming  $\sim 80\%$  simulation time

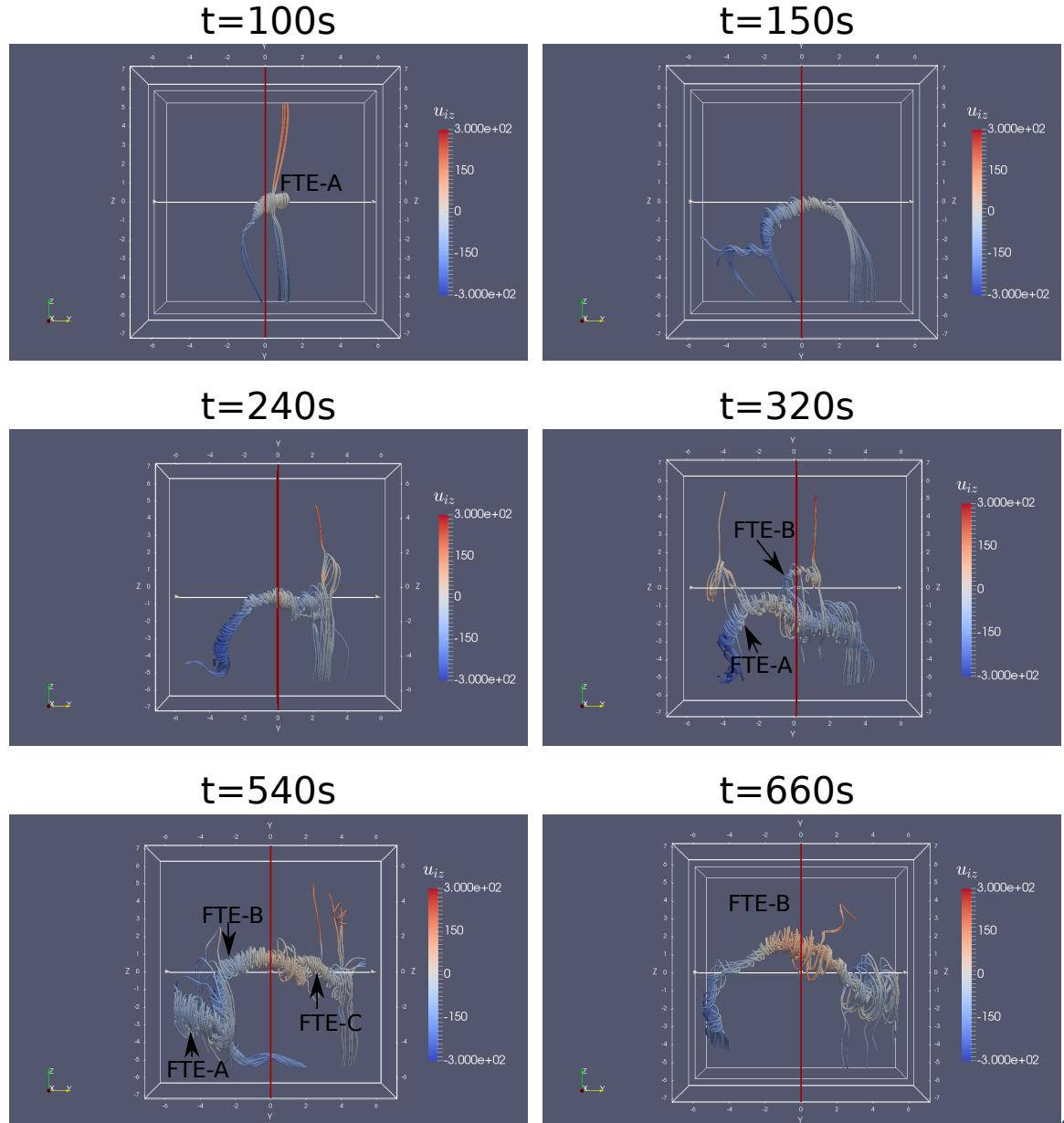
## ➤ $\sim 18000$ core hours modeling 1min



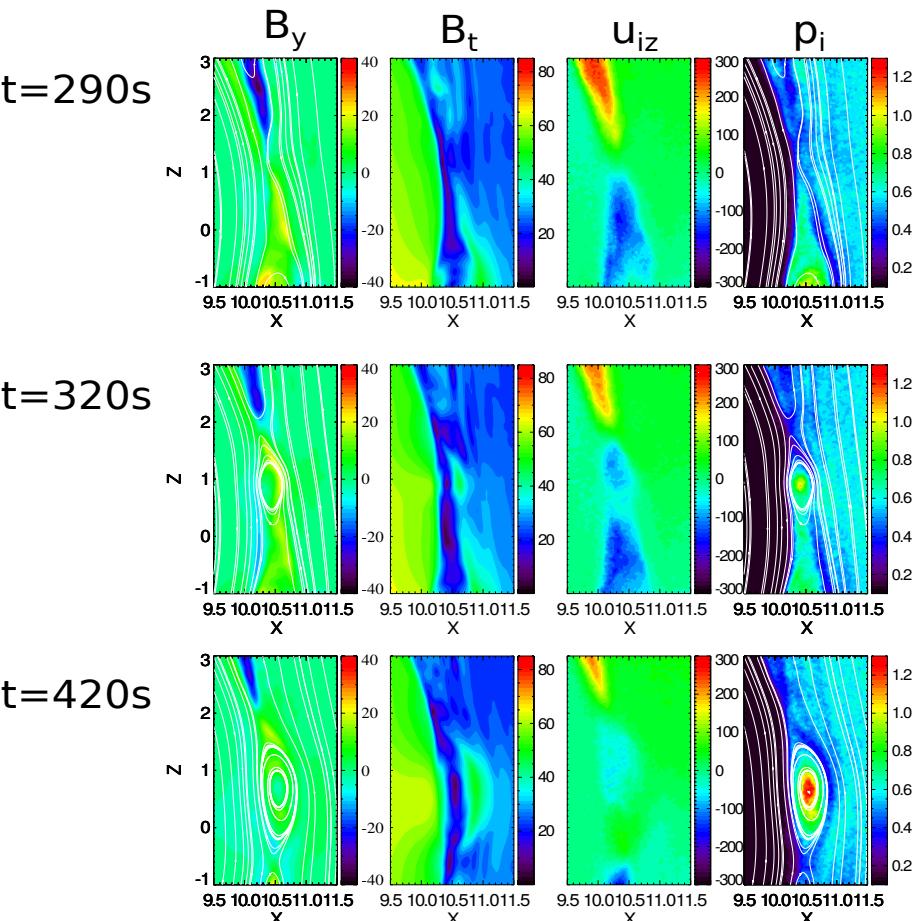
# Evolution of FTEs

## FTEs colored $u_{iz}$

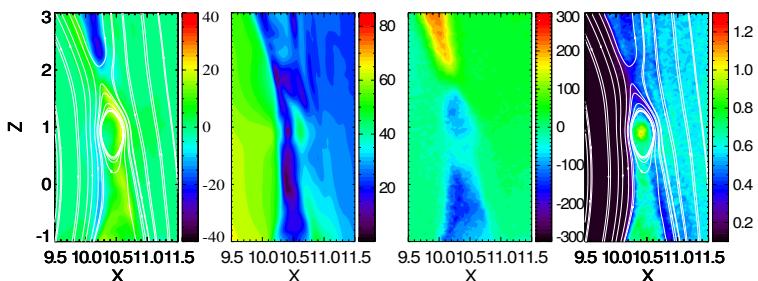
- IMF is purely southward
- FTEs grow in the dawn-dusk direction
- $u_{iz}$  varies along the FTE, and the FTE becomes tilted
- Two FTEs can merge into one



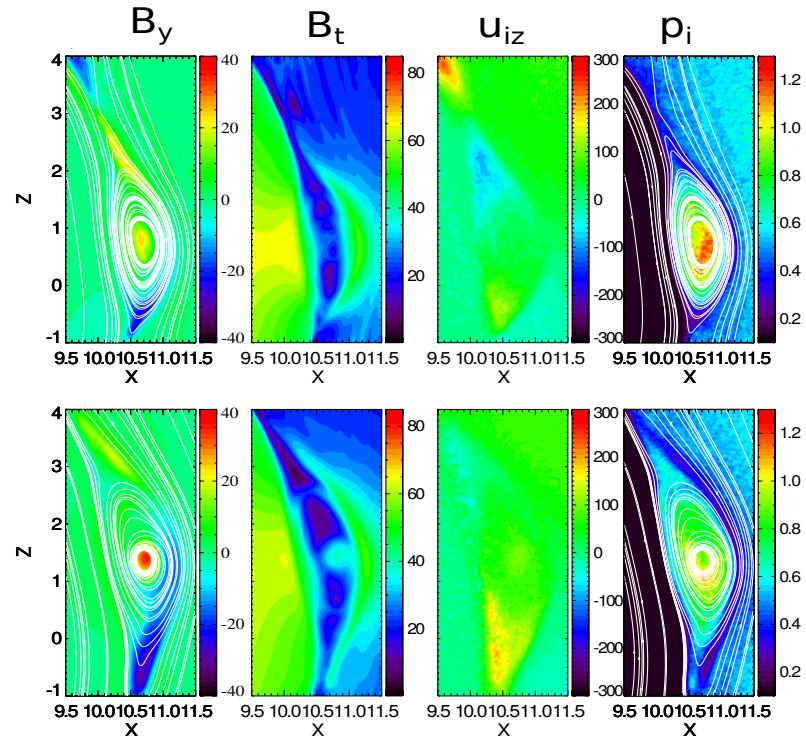
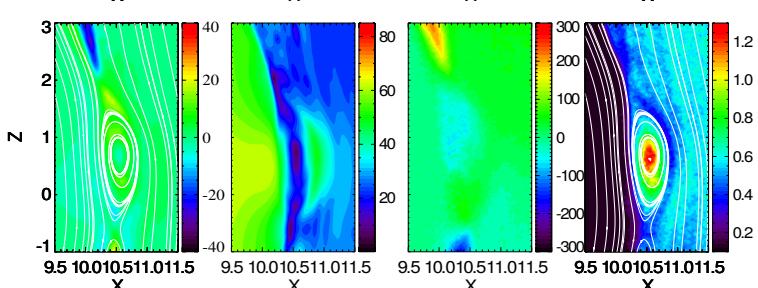
# Evolution of FTEs (2)



$t=540s$

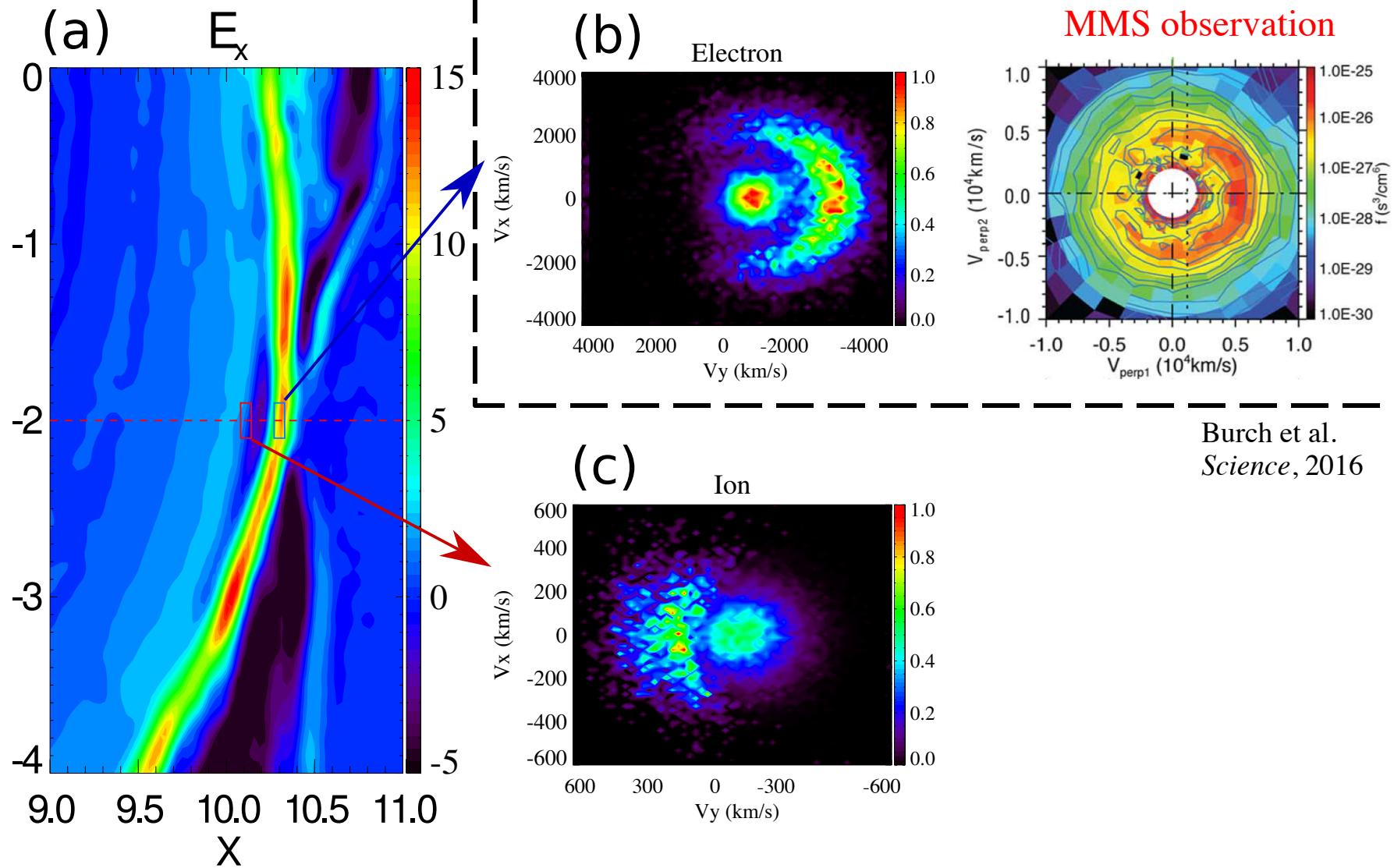


$t=660s$



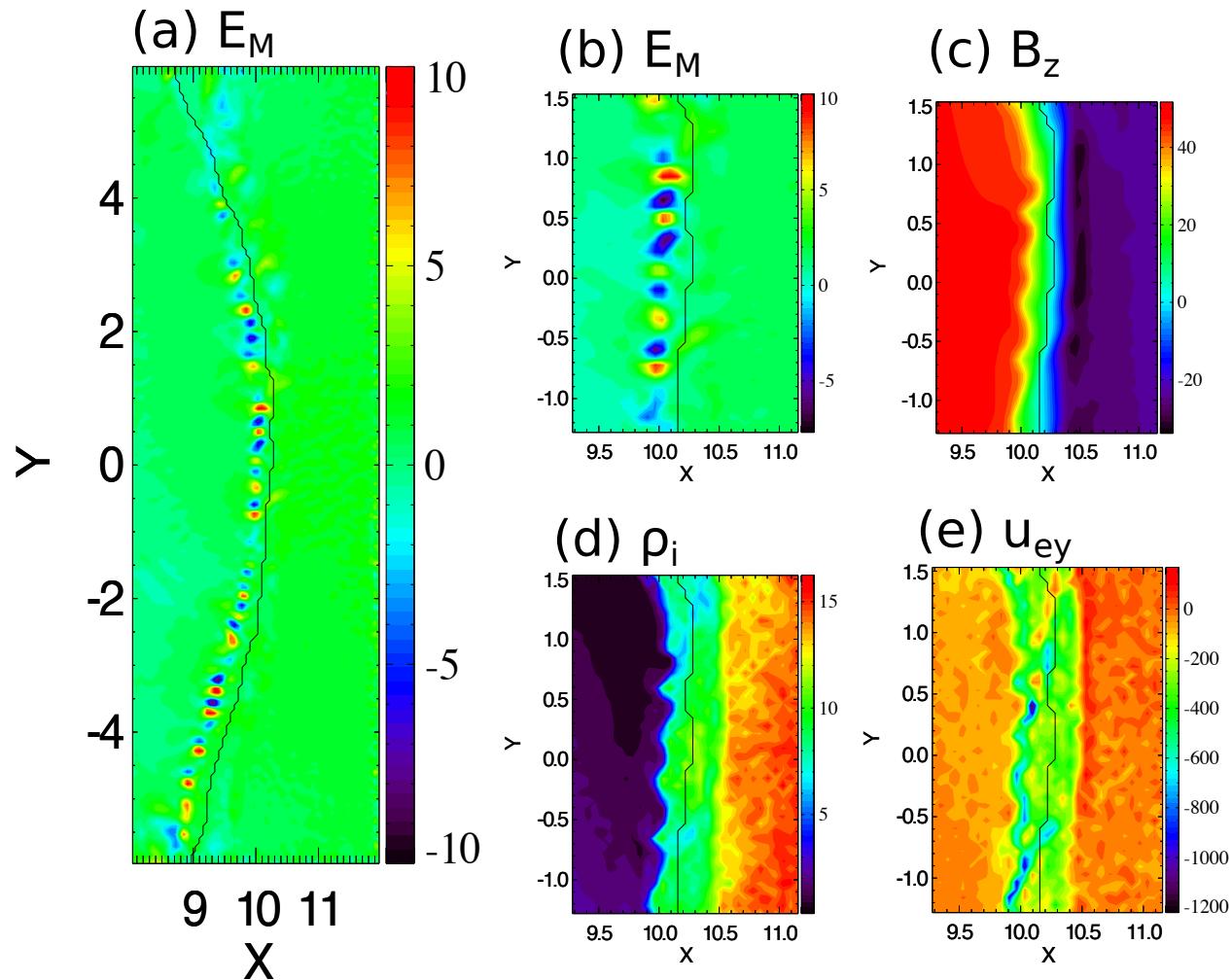
- The FTE is surrounded by two X lines
- The core field gradually increases to a significant value

# Crescent Distributions



# Low Hybrid Drift Instability (LHDI)

- Simulation agrees with MMS observation
  - $E_M \sim 8 \text{ mV/m}$
  - $kr_e \sim 0.4, \lambda \sim 16 r_e$



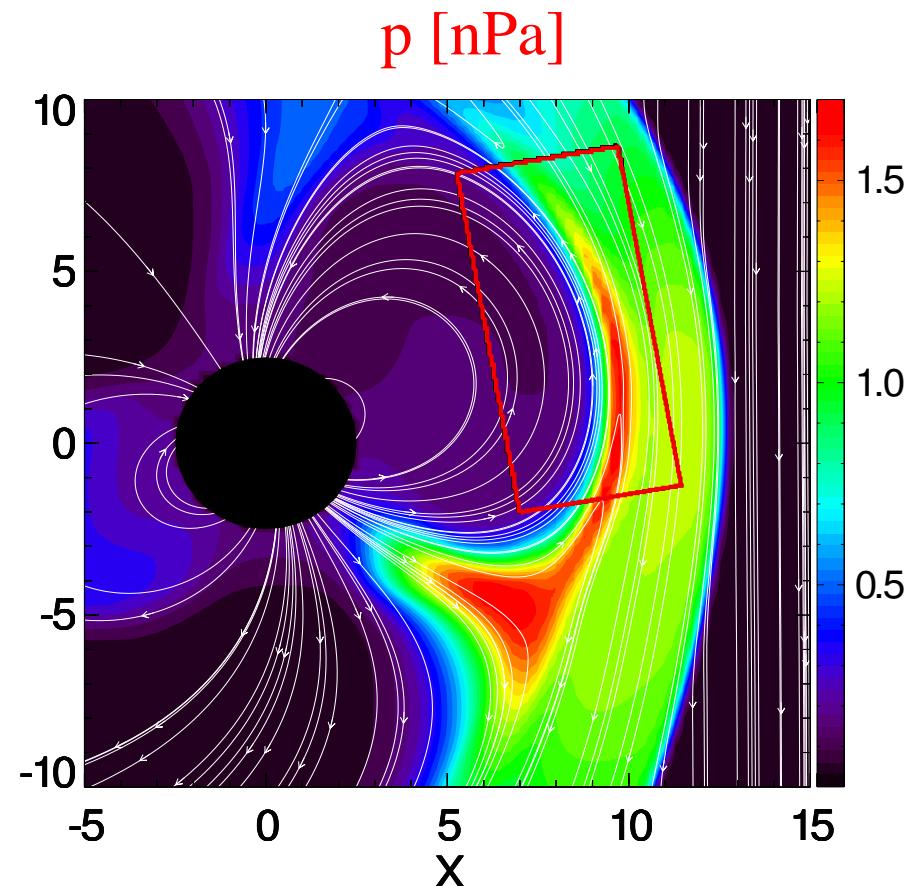
# GEM Dayside Kinetic Challenge Event

➤ The southward IMF event on 2015-11-18  
01:50-03:00 UT

- Solar wind conditions:  
 $\rho = 9.5 \text{ amu/cm}^3$ ,  $U_x = -360 \text{ km/s}$ ,  
 $T_i = T_e = 9 \text{ eV}$ ,  $B = [0, 0, -6] \text{ nT}$
- The dipole is  $27^\circ$  tilted from the Z-axis towards the midnight: rotated PIC region

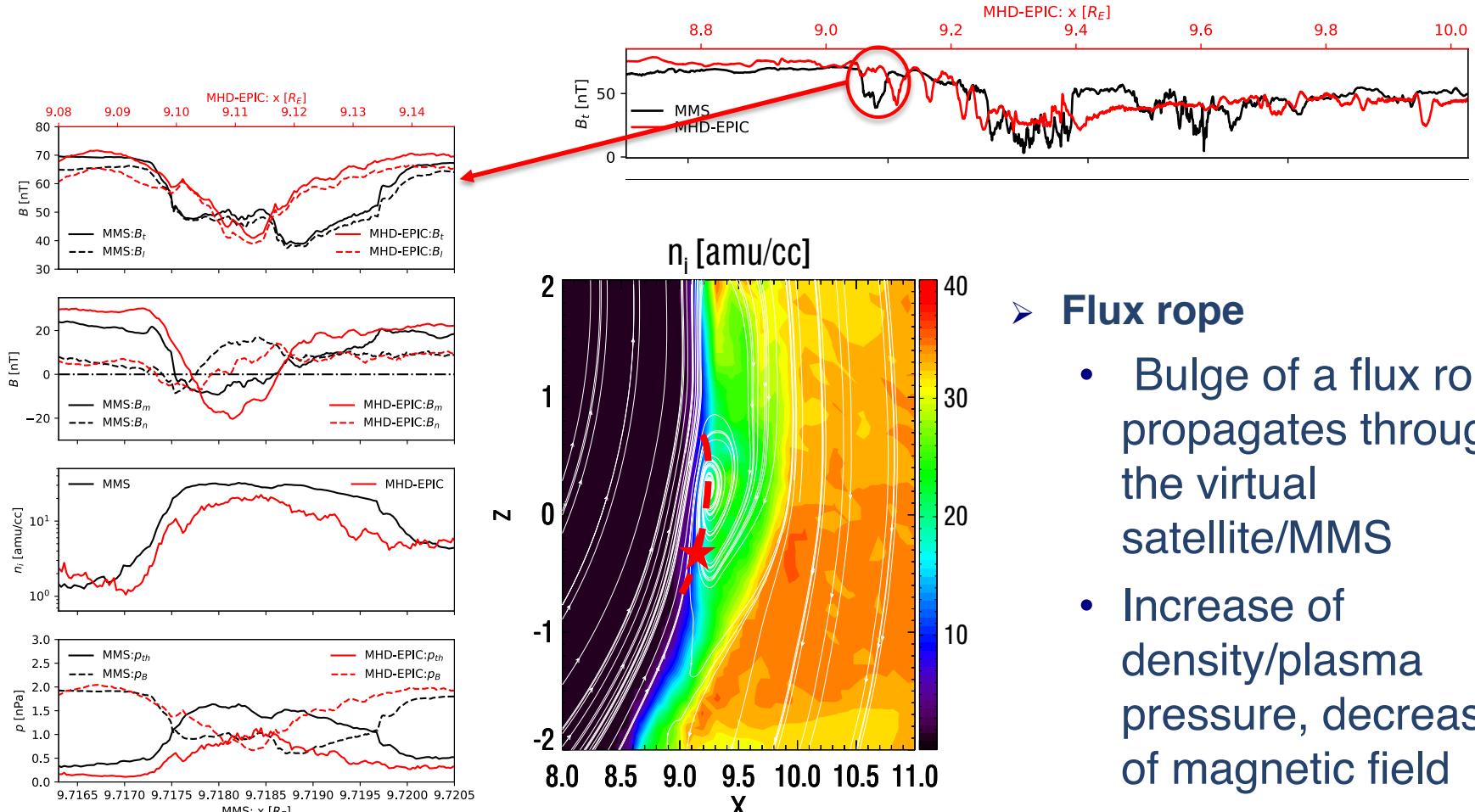
➤ Numerical parameters

- Artificially increase ion kinetic scale by a factor of 16, so  $d_i \sim 1/6 R_E$  in magnetosheath
- $m_i/m_e = 100$
- $\Delta x = 1/25 R_E$ : 4 cells per  $d_i$



Chen et al., 2020, ESS

# MHD-EPIC vs. MMS: FTE-Like Structure

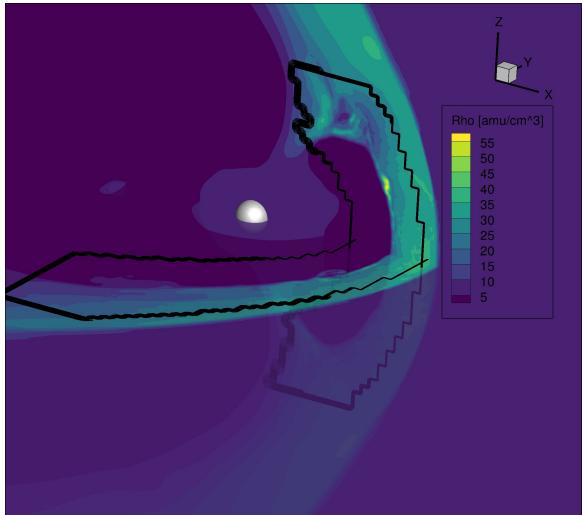


## ➤ Flux rope

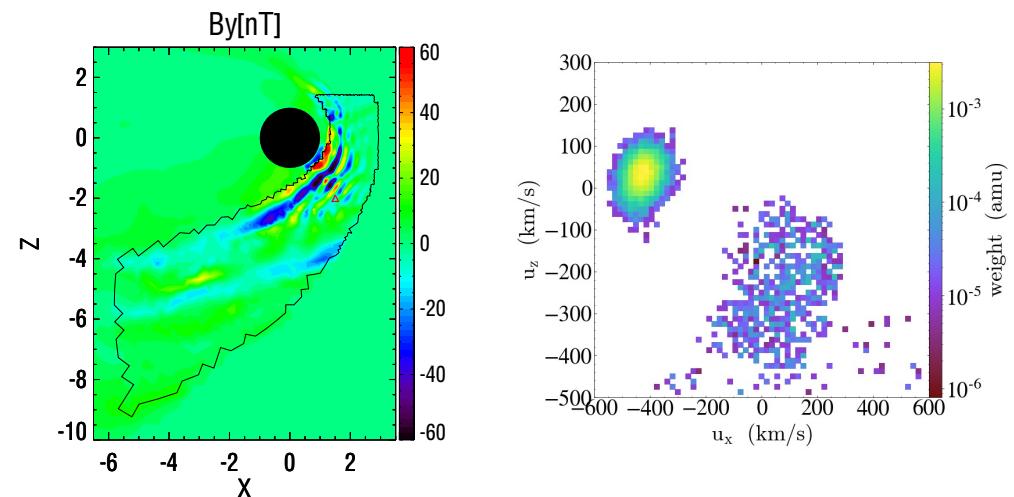
- Bulge of a flux rope propagates through the virtual satellite/MMS
- Increase of density/plasma pressure, decrease of magnetic field
- $B_n$  changes sign

# Simulations with FLEKS

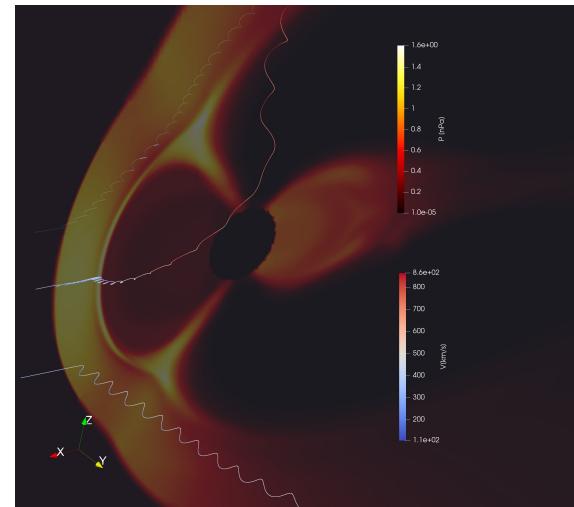
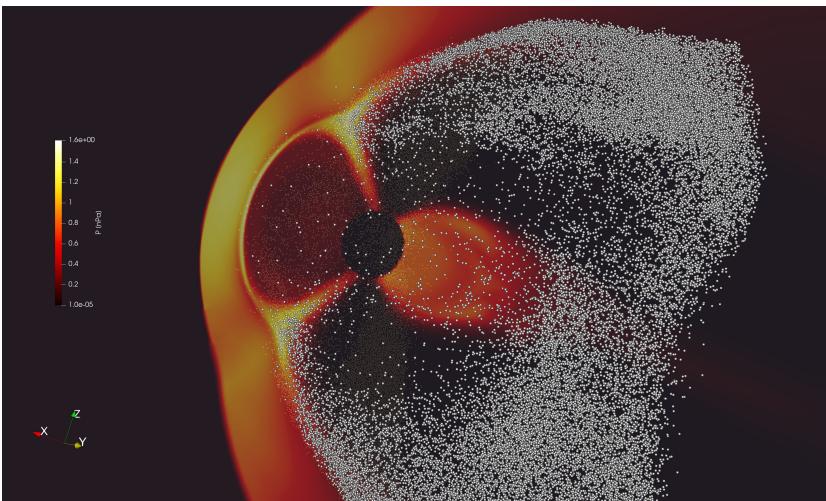
Earth: cusp?



Mercury: foreshock?



FLEKS: newly developed test particle module



# Summary

- **MHD-EPIC couples global MHD model with local PIC**
- **Improvement of the PIC component**
  - Gauss's Law satisfying Energy Conserving Semi-Implicit Method (GL-ECSIM)
    - Adjusts particle positions to satisfy Gauss's law
    - Improves robustness and accuracy
  - FFlexible Exascale Kinetic Simulator (FLEKS)
    - Adaptive PIC region
    - Particle resampling algorithms
- **Applications**
  - Mercury's magnetotail
    - Dawn-dusk asymmetry
  - Earth's dayside magnetopause
    - Asymmetric reconnection
    - Flux transfer events (FTEs)
    - MMS data comparison